









Geophysical Investigation of Asteroids by Spacecraft

UCLA seminar 1 March 2018

A. I. Ermakov¹ (eai@caltech.edu), R. S. Park¹, C. A. Raymond¹, M. T. Bland², M. T. Zuber³, C. T. Russell⁴, R. R. Fu⁵

¹Jet Propulsion Laboratory, California Institute of Technology

²US Geological Survey, Astrogeology Science Center

³Department of the Earth, Atmospheric and Planetary Sciences, Massachusetts Institute of Technology

⁴University of California Los Angeles

⁵Department of Earth and Planetary Sciences, Harvard University.









- > Why do we need to study remnant planetesimals (a.k.a planetary embryos a.k.a protoplanets)?
- > How do we study them?
- Dawn at Vesta
- Dawn at Ceres
- Future studies



Outline





- ➤ Why do we need to study remnant planetesimals (a.k.a planetary embryos a.k.a protoplanets)?
- > How do we study them?
- Dawn at Vesta
- Dawn at Ceres
- Future studies



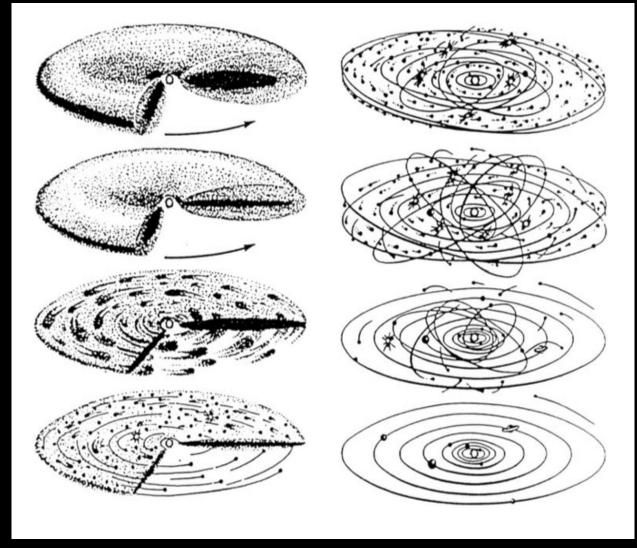




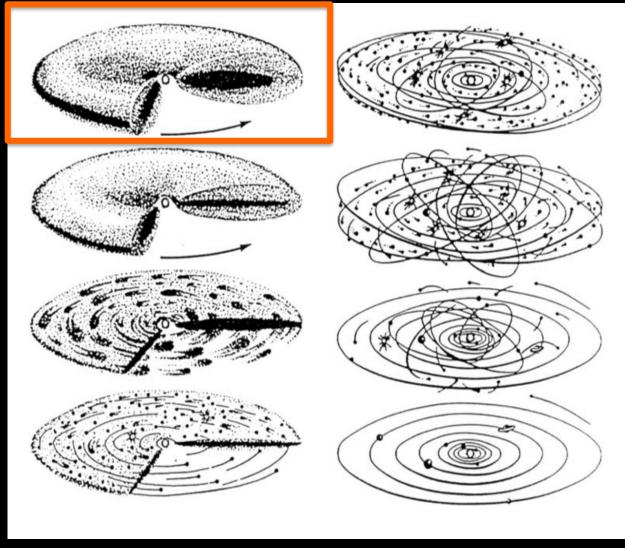


- ➤ Why do we need to study remnant planetesimals (a.k.a planetary embryos a.k.a protoplanets)?
- > How do we study them?
- Dawn at Vesta
- Dawn at Ceres
- Future studies





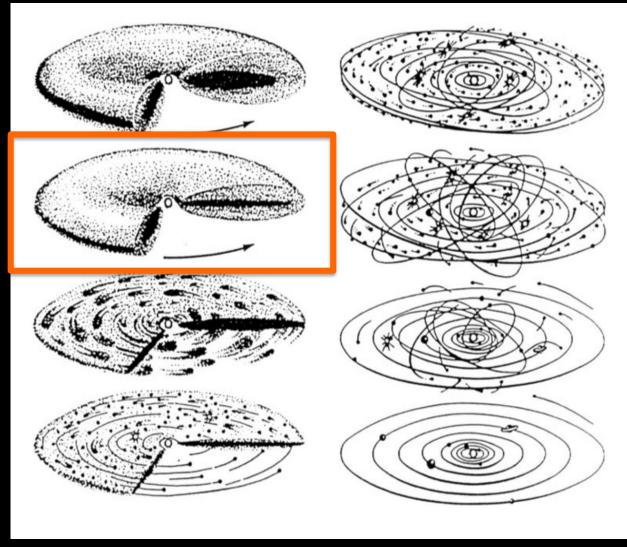




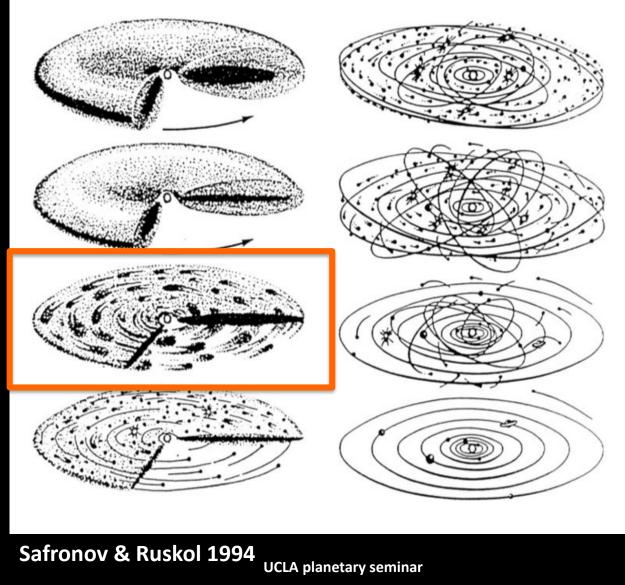




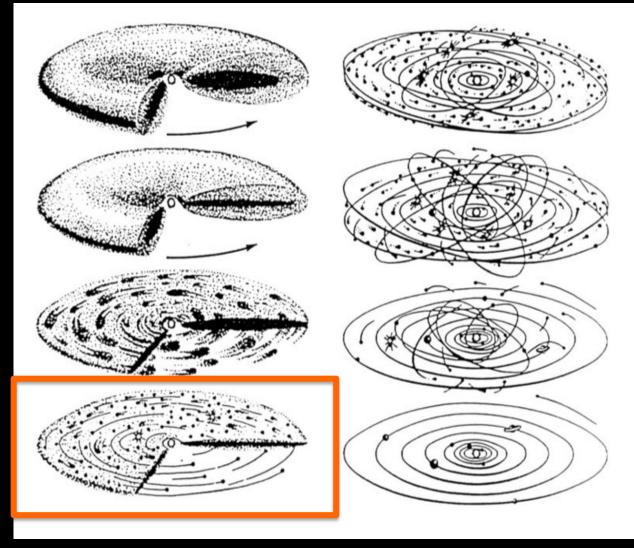


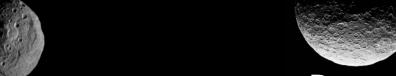






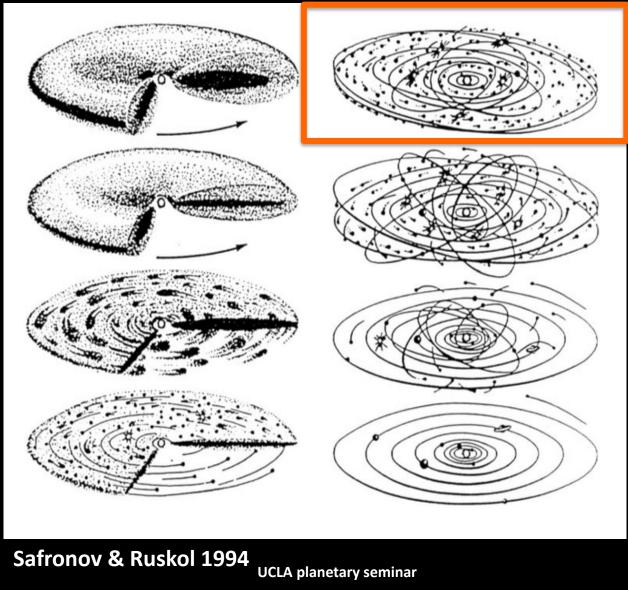




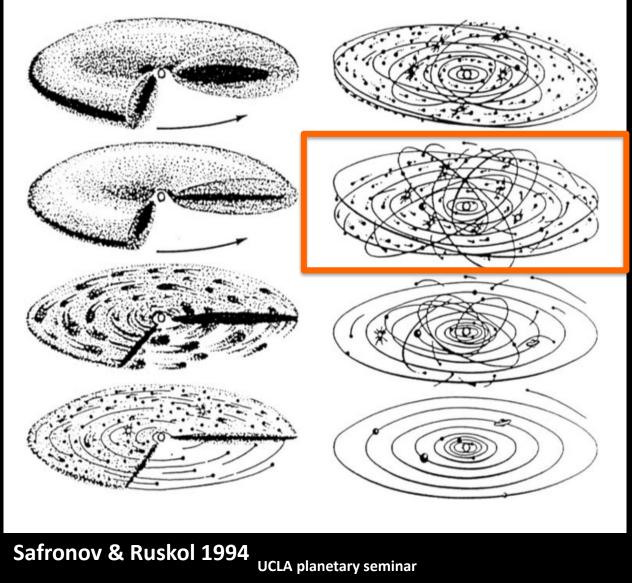




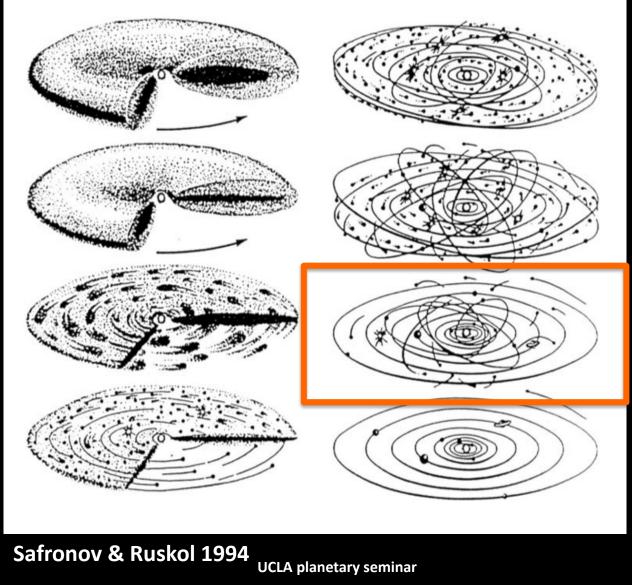
Run-away growth



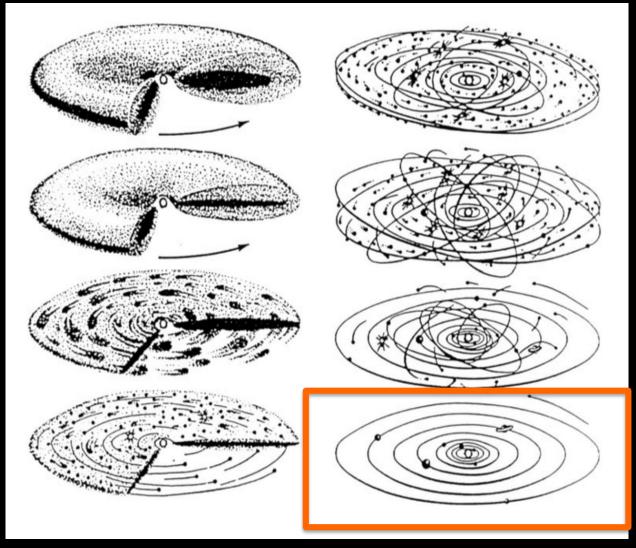








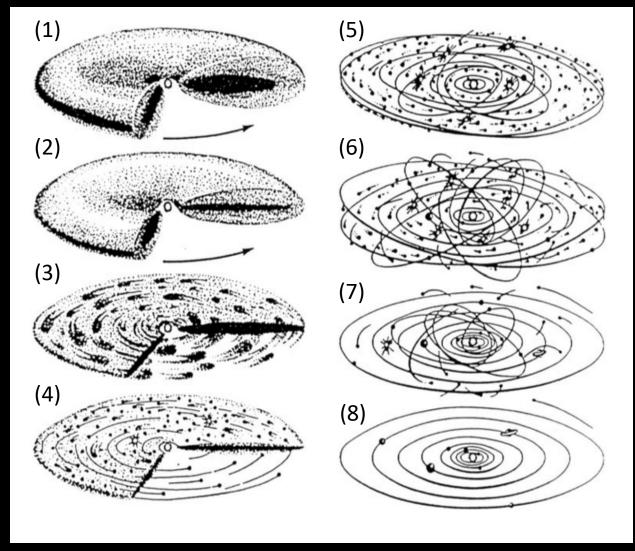






Planet formation

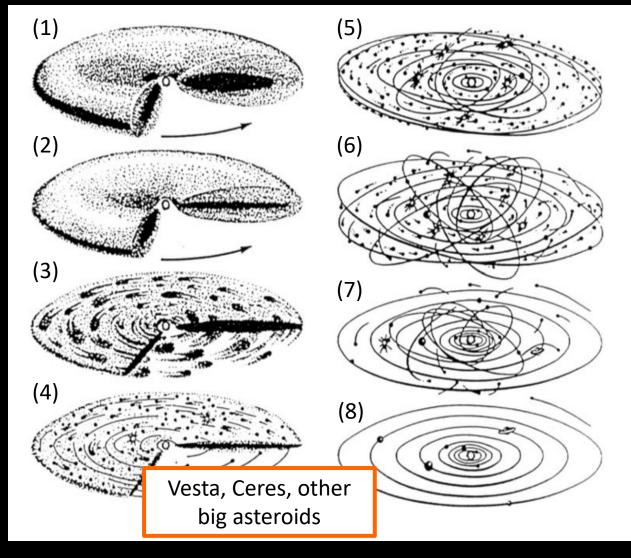
- 1 Formation of a nebula disk
- 2 Settling to midplane
- 3 Dust coagulation
- 4 Orderly growth
- (5) Run-away growth
- **6** Gas dispersal
- **7** Late-state mergers
- 8 Present state



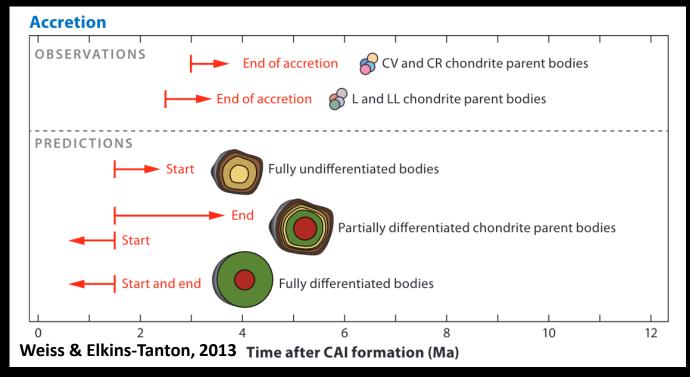


Planet formation

- 1 Formation of a nebula disk
- ② Settling to midplane
- 3 Dust coagulation
- 4 Orderly growth
- (5) Run-away growth
- **6** Gas dispersal
- **7** Late-state mergers
- 8 Present state







- What was the differentiation state of planetesimals?
 - Differentiated or undifferentiated?
 - How much water?
- What can interior structure tell us about the accretion process?
 - Fast or slow
 - Early or late



How do we study a planetary interior with gravity and topography?

- > We study the interior but looking at its response to various forcings such as:
 - Rotation
 - Surface loads
 - Subsurface loads











- In hydrostatic equilibrium
 - Surfaces of constant density, pressure and potential coincide
 - No shear stresses

























$$\rho = \rho(r)$$
, ω







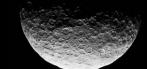


Hydrostatic equilibrium

$$\rho = \rho(r)$$
, ω



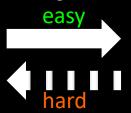








$$\rho = \rho(r)$$
, ω

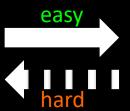






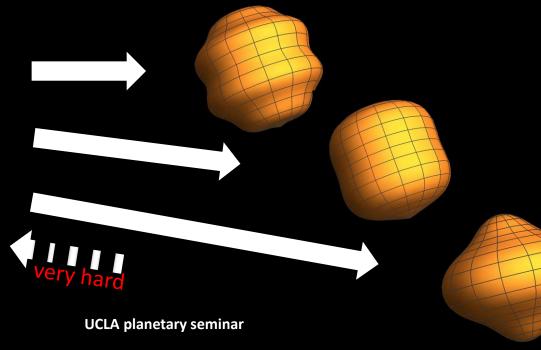
Hydrostatic equilibrium

$$\rho = \rho(r)$$
, ω





$$\rho = \rho(r)$$
, ω











Shape

Spherical Harmonics

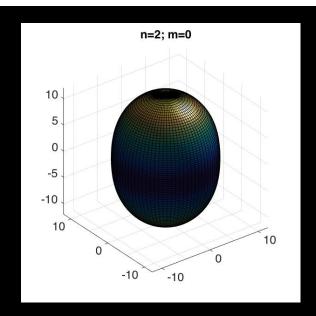
Gravitational potential

$$U(r,\phi,\lambda) = \frac{GM}{R} \left[1 + \sum_{n=2}^{\infty} \sum_{m=0}^{n} \left(\frac{R_0}{r} \right)^n \left(C_{nm} \cos(m\lambda) + S_{nm} \sin(m\lambda) \right) P_{nm} \left(\sin \phi \right) \right]$$

Power Spectral Density

$$S_n^{gg} = \mathop{a}\limits_{m=0}^n \frac{C_{nm}^2 + S_{nm}^2}{2n+1}$$
 gravity

$$S_n^{tt} = \mathop{a}\limits_{m=0}^n \frac{A_{nm}^2 + B_{nm}^2}{2n+1}$$
 topography











Spherical Harmonics

Shape

Gravitational potential

$$U(r,\phi,\lambda) = \frac{GM}{R} \left[1 + \sum_{n=2}^{\infty} \sum_{m=0}^{n} \left(\frac{R_0}{r} \right)^n \left(C_{nm} \cos(m\lambda) + S_{nm} \sin(m\lambda) \right) P_{nm} \left(\sin \phi \right) \right]$$

U – gravitational potential

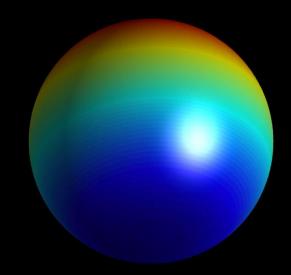
 φ – latitude

 λ – longitude

r – radial distance

n – degree

m – order













Shape

Spherical Harmonics

$$r(f, I) = R_0 \stackrel{\text{\'e}}{\overset{\text{\'e}}}{\overset{\text{\'e}}{\overset{\text{\'e}}{\overset{\text{\'e}}}{\overset{\text{\'e}}{\overset{\text{\'e}}}{\overset{\text{\'e}}{\overset{\text{\'e}}{\overset{\text{\'e}}}{\overset{\text{\'e}}{\overset{\text{\'e}}}{\overset{\text{\'e}}{\overset{\text{\'e}}}{\overset{\text{\'e}}{\overset{\text{\'e}}}{\overset{\text{\'e}}}{\overset{\text{\'e}}}}}}}}}}}}}} (A_{nm}\cos(mI) + B_{nm}\sin(mI)) + B_{nm}\sin(mI)})P_{nm}\sin(mI)} \stackrel{\text{\'e}}{\overset{\text{\'e}}}{\overset{\text{\'e}}}}}}} \stackrel{\text{\'e}}{\overset{\text{\'e}}}}}{\overset{\text{\'e}}{\overset{\text{\'e}}}}}} (A_{nm}\cos(mI) + B_{nm}\sin(mI)}) + B_{nm}\sin(mI)) + B_{nm}\sin(mI)} \stackrel{\text{\'e}}{\overset{\text{\'e}}}}}}{\overset{\text{\'e}}}{\overset{\text{\'e}}}}}}} (A_{nm}\cos(mI) + B_{nm}\cos(mI)}) + B_{nm}\sin(mI)}) + B_{nm}\sin(mI)} \stackrel{\text{\'e}}{\overset{\text{\'e}}}}}}{\overset{\text{\'e}}}}} (A_{nm}\cos(mI) + B_{nm}\cos(mI)}) + B_{nm}\sin(mI)}) + B_{nm}\sin(mI)} \stackrel{\text{\'e}}{\overset{\text{\'e}}}}}{\overset{\text{\'e}}}} (A_{nm}\cos(mI) + B_{nm}\cos(mI)}) + B_{nm}\sin(mI)} \stackrel{\text{\'e}}{\overset{\text{\'e}}}}} (A_{nm}\cos(mI) + B_{nm}\cos(mI)}) + B_{nm}\sin(mI)) + B_{nm}\sin(mI)} \stackrel{\text{\'e}}{\overset{\text{\'e}}}} (A_{nm}\cos(mI) + B_{nm}\sin(mI)) + B_{nm}\sin(mI)} \stackrel{\text{\'e}}{\overset{\text{\'e}}}} (A_{nm}\cos(mI) + B_{nm}\sin(mI)}) + B_{nm}\sin(mI)} \stackrel{\text{\'e}}{\overset{\text{\'e}}} (A_{nm}\cos(mI) + B_{nm}\sin(mI)}) + B_{nm}\sin(mI)} \stackrel{\text{\'e}}{\overset{\text{\'e}}} (A_{nm}\cos(mI) + B_{nm}\sin(mI)}) + B_{nm}\sin(mI)} \stackrel{\text{\'e}} (A_{nm}\cos(mI) + B_{nm}\sin(mI)) + B_$$

Gravitational potential

$$U(r,\phi,\lambda) = \frac{GM}{R} \left[1 + \sum_{n=2}^{\infty} \sum_{m=0}^{n} \left(\frac{R_0}{r} \right)^n \left(C_{nm} \cos(m\lambda) + S_{nm} \sin(m\lambda) \right) P_{nm} \left(\sin \phi \right) \right]$$

U – gravitational potential

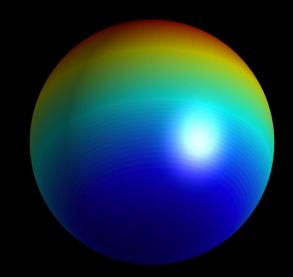
 φ – latitude

 λ – longitude

r – radial distance

n – degree

m – order













Z_n - gravity-topography admittance

$$Z_n = \frac{\text{gravity}}{\text{topography}}$$
 for a given wavelength









Admittance

Z_n - gravity-topography admittance

$$Z_n = \frac{\text{gravity}}{\text{topography}}$$
 for a given wavelength

Linear two-layer hydrostatic model

$$Z_n = \frac{GM}{R^3} \frac{3(n+1)}{2n+1} \frac{\Gamma_{crust}}{\Gamma_{mean}}$$









Admittance

Z_n - gravity-topography admittance

$$Z_n = \frac{\text{gravity}}{\text{topography}}$$
 for a given wavelength

Linear two-layer hydrostatic model

$$Z_n = \frac{GM}{R^3} \frac{3(n+1)}{2n+1} \frac{\Gamma_{crust}}{\Gamma_{mean}}$$

> Linear two-layer isostatic model

$$Z_{n} = \frac{GM}{R^{3}} \frac{3(n+1)}{2n+1} \frac{\Gamma_{crust}}{\Gamma_{mean}} \hat{\mathbf{e}}^{1} - \mathbf{e}^{1} - \frac{D_{comp}}{R} \ddot{\mathbf{e}}^{n} \dot{\mathbf{u}}$$











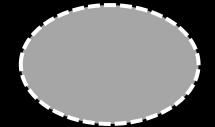
Gravity anomalies

Free-air anomaly

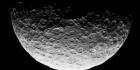
$$\sigma_{\mathsf{FA}} = \sigma_{\mathsf{obs}} - \sigma_{\mathsf{model}}$$

$$\sigma_{
m model}$$
 =

gravity of hydrostatic figure











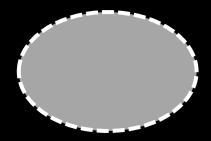


Free-air anomaly

$$\sigma_{\mathsf{FA}} = \sigma_{\mathsf{obs}} - \sigma_{\mathsf{model}}$$

$$\sigma_{
m model}$$
 =

gravity of hydrostatic figure



Bouguer anomaly

$$\sigma_{\mathsf{BA}} = \sigma_{\mathsf{obs}} - \sigma_{\mathsf{model}}$$

$$\sigma_{\mathsf{model}} =$$

gravity of shape assuming p



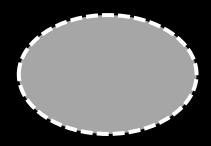


Gravity anomalies

Free-air anomaly

$$\sigma_{\mathsf{FA}} = \sigma_{\mathsf{obs}} - \sigma_{\mathsf{model}}$$

$$\sigma_{\text{model}} =$$
 gravity of hydrostatic figure



Bouguer anomaly

$$\sigma_{\mathsf{BA}} = \sigma_{\mathsf{obs}} - \sigma_{\mathsf{model}}$$

$$\sigma_{\mathsf{model}} =$$

gravity of shape assuming ρ

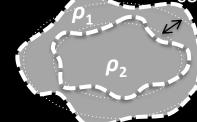
Isostatic anomaly

$$\sigma_{\mathsf{IA}} = \sigma_{\mathsf{obs}} - \sigma_{\mathsf{model}}$$

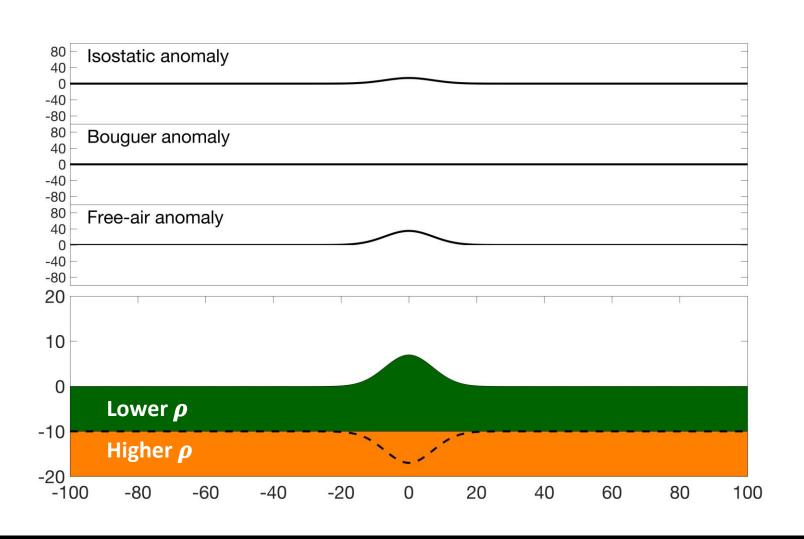
h – depth of

compensation

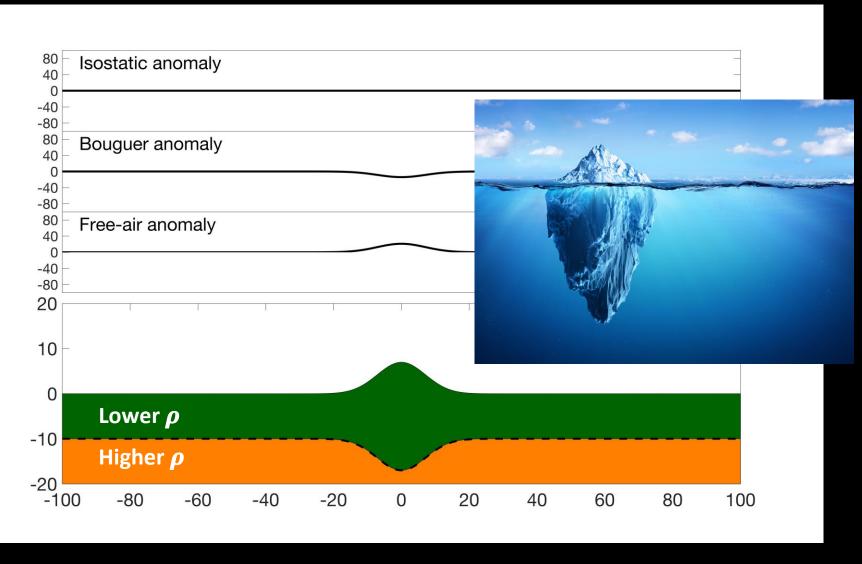
$$\sigma_{\text{model}} =$$
 gravity assuming isostasy for ρ_1, ρ_2, h



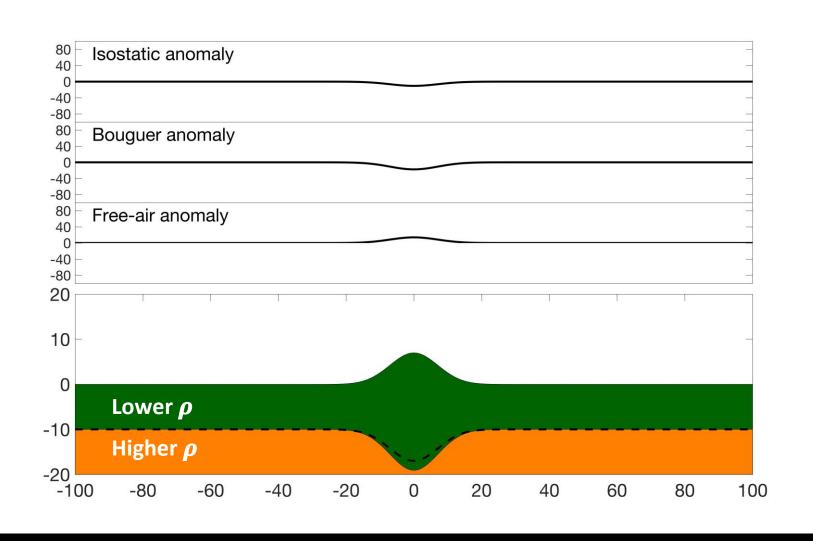




Example: compensated topography















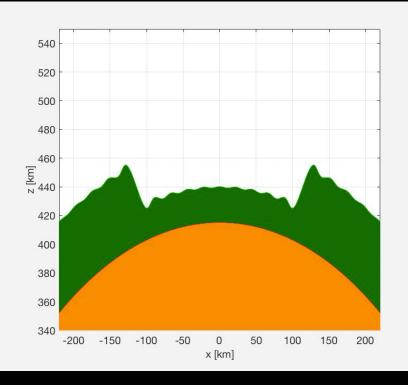


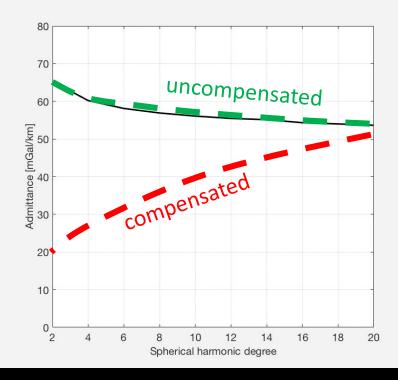
Isostatic compensation

Example of a spherical cap (depression) relaxation

Interface evolution

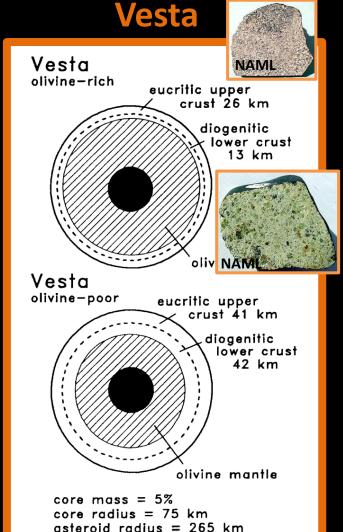
Admittance evolution

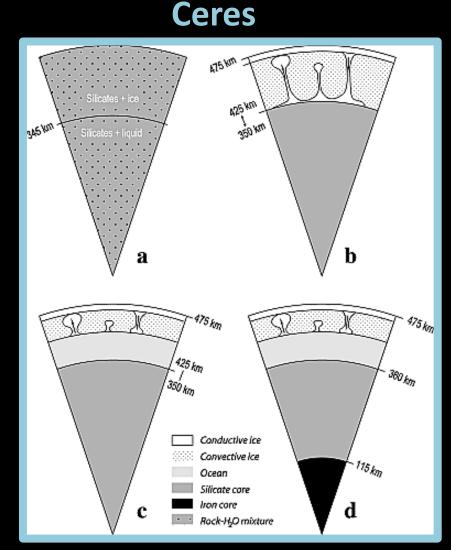






What did we know before Dawn?

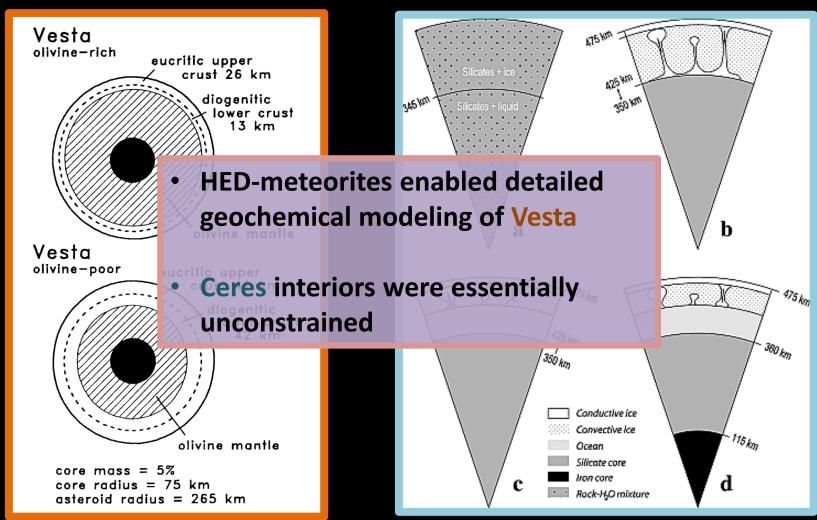








What did we know before Dawn?



Ruzicka et al., 1997

McCord and Sotin, 2005





Dawn geophysical data

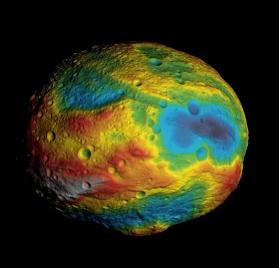
- Shape model
 - Stereophotogrammetry (SPG) from DLR
 - Stereophotoclinometry (SPC) from JPL
 - Mutually consistent with the accuracy much better than the spatial resolution of gravity field
- Gravity field
 - Accurate up to n = 18 ($\lambda = 93$ km) for Vesta (Konopliv et al., 2014)
 - Accurate up to n = 17 ($\lambda = 174$ km) for Ceres (Konopliv et al., 2017)
- Assumptions we have to make:
 - Multilayer model with uniform density layers
 - Range of core densities for Vesta
 - Range of crustal densities from HEDs for Vesta
 - Can't really assume anything for Ceres

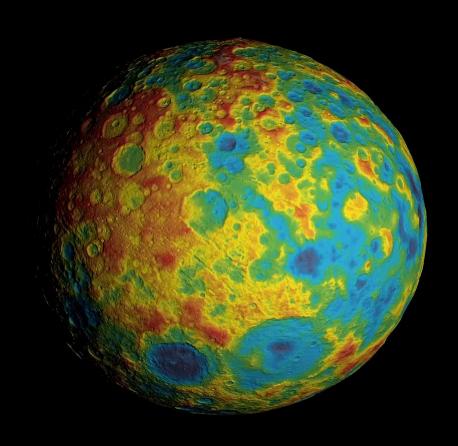












Gaskell, 2012

Park et al., 2016

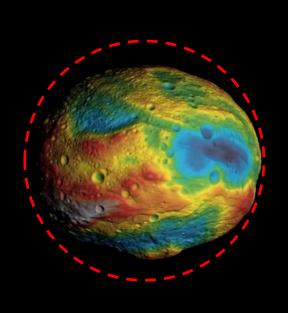
UCLA planetary seminar

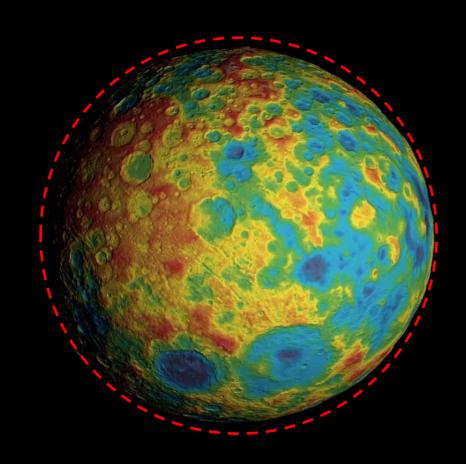












Gaskell, 2012

Park et al., 2016

UCLA planetary seminar



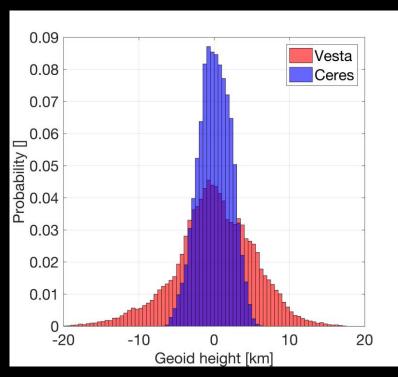


Shape statistics

Parameter	Vesta	Ceres
Equatorial flattening	0.0262	0.0043
Geoidal height range (km)	37.9	13.2
Geoidal height RMS (km)	5.2	2.1

- Ceres is closer to hydrostatic equilibrium than Vesta
- Smoother topography at Ceres

Hypsograms of Vesta and Ceres



*Hypsogram is a fancy word for the "histogram of elevations"

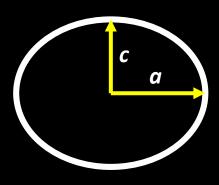






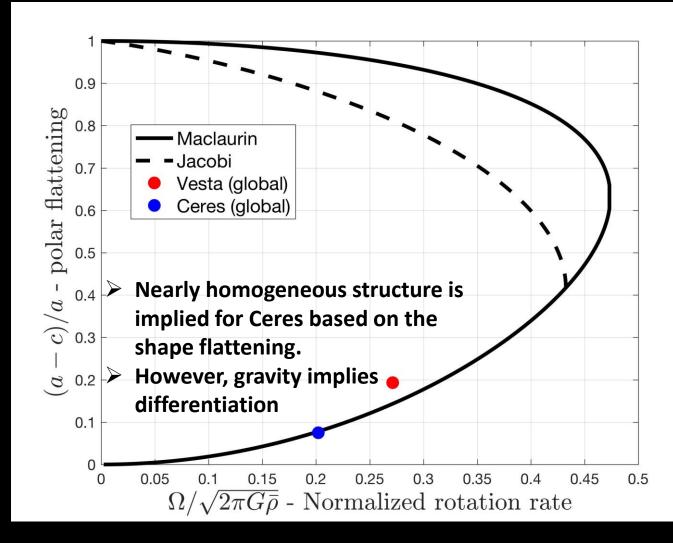


Flattening vs rotation rate



homogeneous more oblate

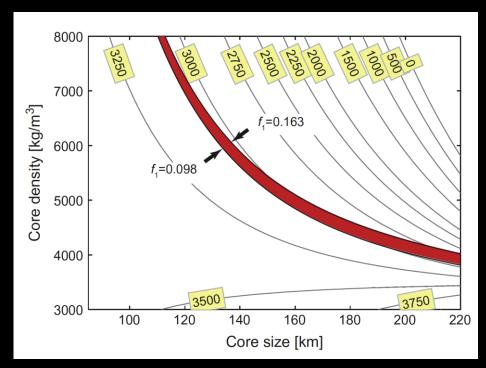
differentiated less oblate







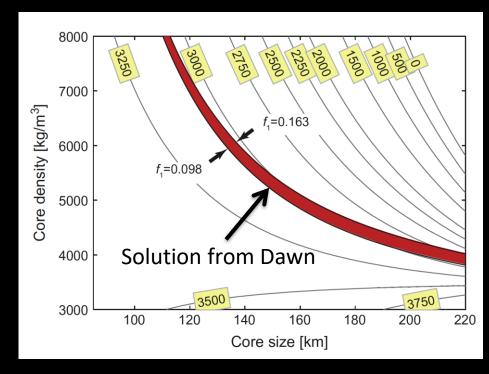
- Vesta is not presently in hydrostatic equilibrium
- No unique solution only from gravity/topography, need an extra constraint
- Geochemically motivated 3layer interior structure (Ruzicka et al., 1997)
- Densities constrained by the Howardite-Eucrite-Diogenite (HED) meteorites







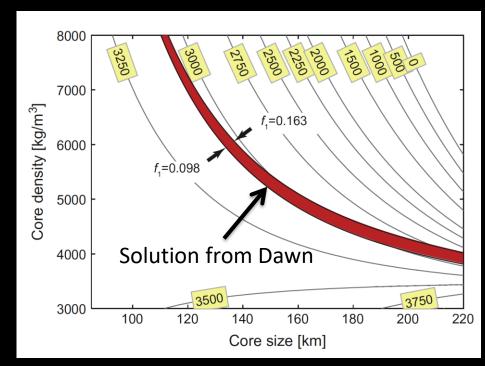
- Vesta is not presently in hydrostatic equilibrium
- No unique solution only from gravity/topography, need an extra constraint
- Geochemically motivated 3layer interior structure (Ruzicka et al., 1997)
- Densities constrained by the Howardite-Eucrite-Diogenite (HED) meteorites





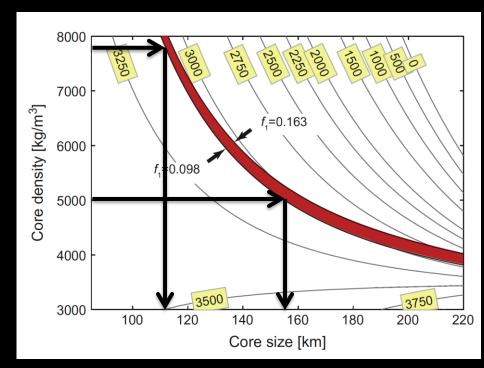


- Vesta is not presently in hydrostatic equilibrium
- No unique solution only from gravity/topography, need an extra constraint
- Geochemically motivated 3layer interior structure (Ruzicka et al., 1997)
- Densities constrained by the Howardite-Eucrite-Diogenite (HED) meteorites





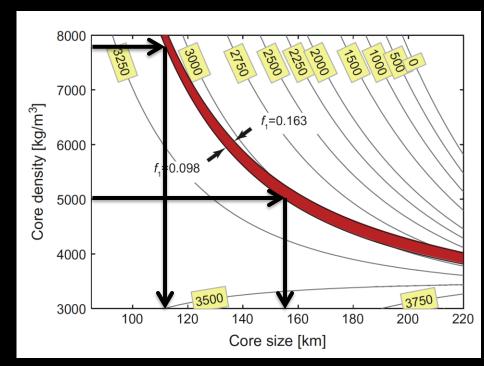
- Vesta is not presently in hydrostatic equilibrium
- No unique solution only from gravity/topography, need an extra constraint
- Geochemically motivated 3layer interior structure (Ruzicka et al., 1997)
- Densities constrained by the Howardite-Eucrite-Diogenite (HED) meteorites





- Vesta is not presently in hydrostatic equilibrium
- No unique solution only from gravity/topography, need an extra constraint
- Geochemically motivated 3layer interior structure (Ruzicka et al., 1997)
- Densities constrained by the Howardite-Eucrite-Diogenite (HED) meteorites

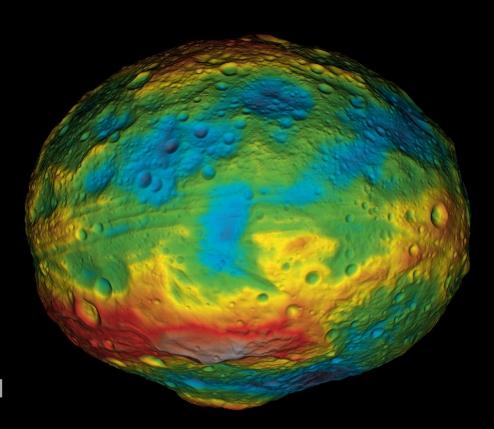
Contours are mantle density [kg/m³]



Core radius of 110 to 155 km Ermakov et al., 2014

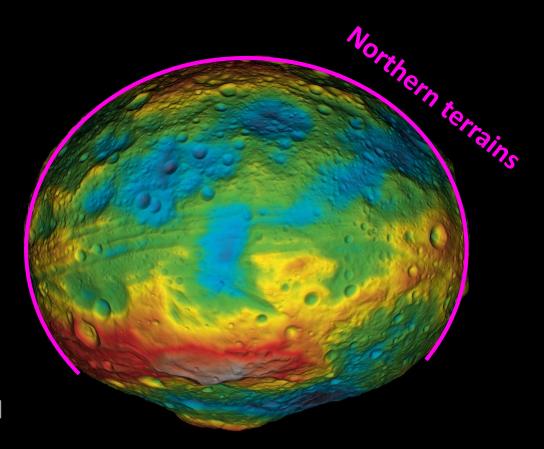


- Vesta was likely close to hydrostatic equilibrium in its early history (Fu et al., 2014).
- Vesta's northern terrains likely reflect its pre-impact equilibrium shape.
- Major impact occurred when Vesta was effectively nonrelaxing leading to uncompensated Rheasilvia and Veneneia basins.



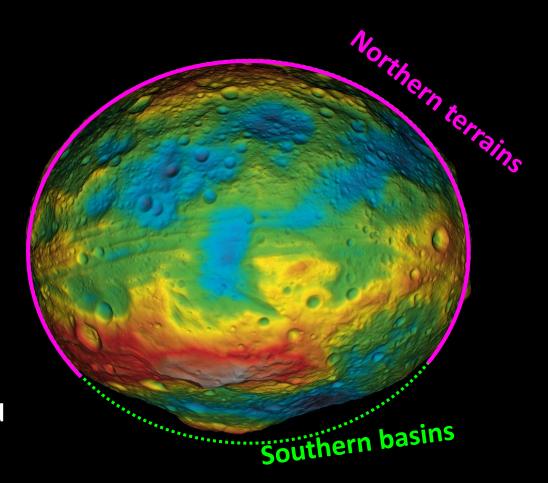


- Vesta was likely close to hydrostatic equilibrium in its early history (Fu et al., 2014).
- Vesta's northern terrains likely reflect its pre-impact equilibrium shape.
- Major impact occurred when Vesta was effectively nonrelaxing leading to uncompensated Rheasilvia and Veneneia basins.

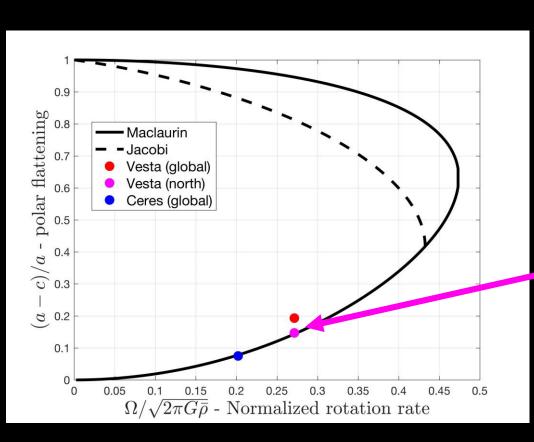


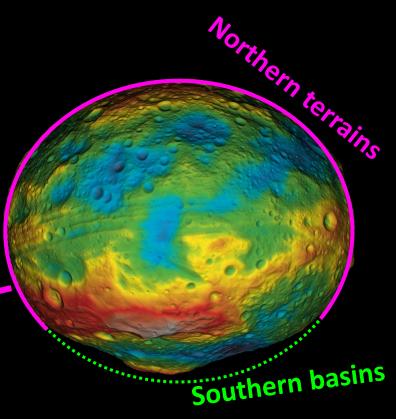


- Vesta was likely close to hydrostatic equilibrium in its early history (Fu et al., 2014).
- Vesta's northern terrains likely reflect its pre-impact equilibrium shape.
- Major impact occurred when Vesta was effectively nonrelaxing leading to uncompensated Rheasilvia and Veneneia basins.

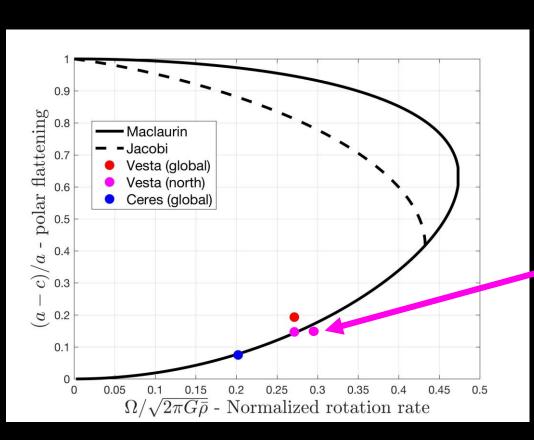


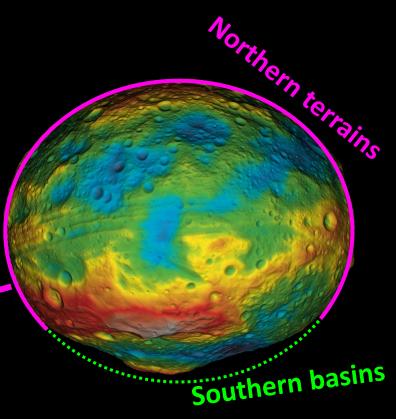
Early efficient viscous relaxation of Vesta







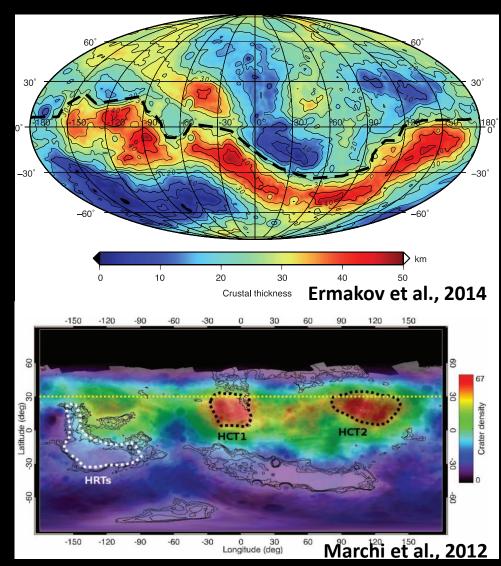






Crustal thickness inversion show a belt of thicker crust around the Southern Basins

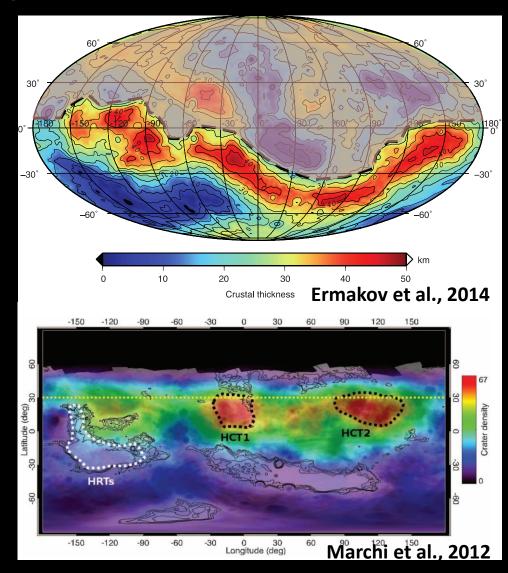
Crater counting reveals that the northern Vesta terrains are old (>3Gy)



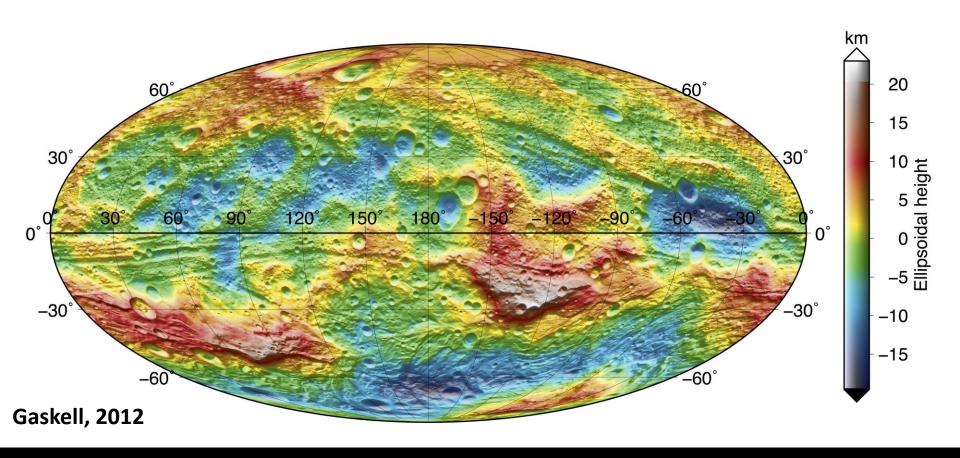


Crustal thickness inversion show a belt of thicker crust around the Southern Basins

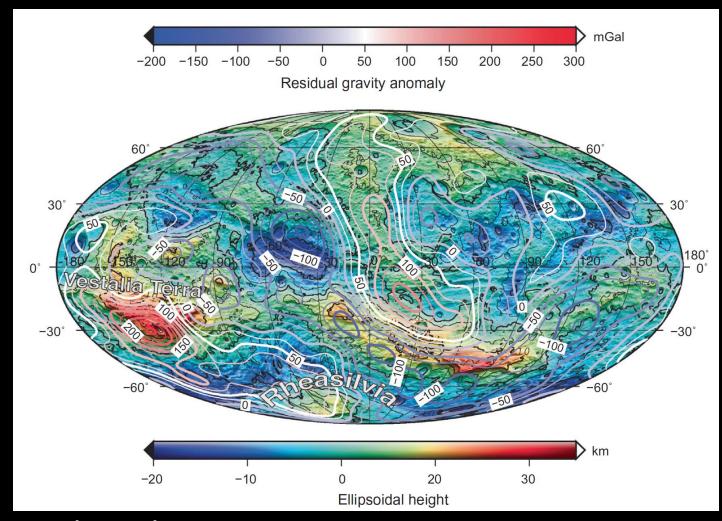
Crater counting reveals that the northern Vesta terrains are old (>3Gy)











Summary on Vesta

- Formed early (< 5 My after CAI)</p>
- Once hot and hydrostatic, Vesta is no longer either
- Differentiated interior
- Most of topography acquired when Vesta was already cool => uncompensated topography
- Combination of gravity/topography data with meteoritic geochemistry data provides constraints on the internal structure





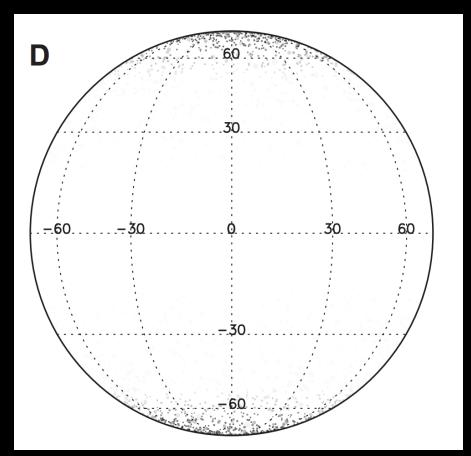






Ceres Expectations

- Bland et al., 2013 predicted that craters on Ceres would quickly relax in an icedominated shell
 - Equatorial warmer craters would relax faster than colder polar craters
- Bland et al., 2016 did not find evidence for such relaxation pattern
 - No latitude dependence of crater depth



Bland, 2013









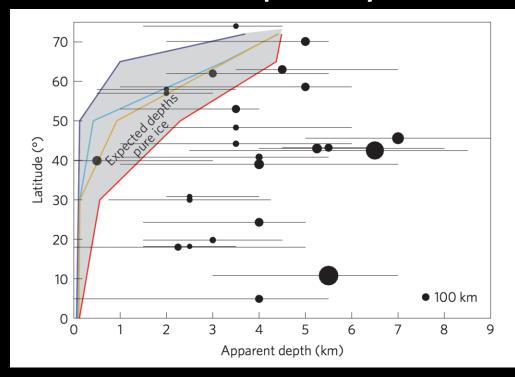


Ceres observation

- Bland et al., 2013 predicted that craters on Ceres would quickly relax in an icedominated shell
 - Equatorial warmer craters would relax faster than colder polar craters
- Bland et al., 2016 did not find evidence for such relaxation pattern

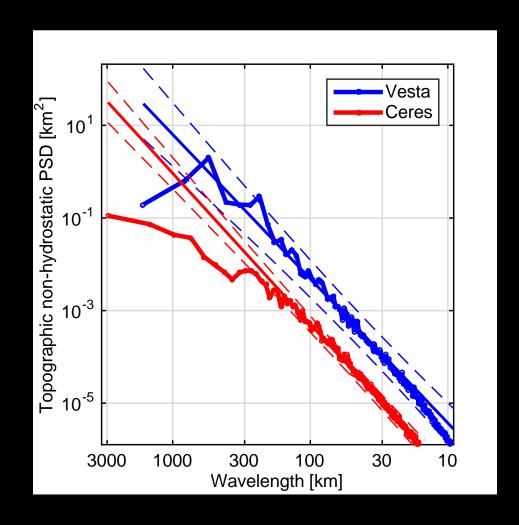
No latitude dependence of crater depth

Crater depth study



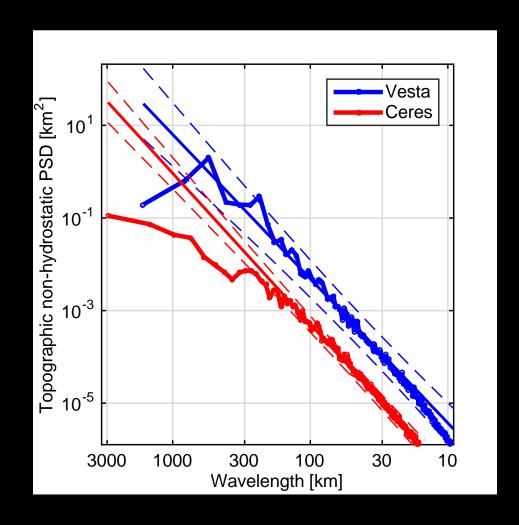


- More general approach: <u>study topography power</u> <u>spectrum</u>
- Power spectra for Vesta closely fits with the power law to the lowest degrees (λ < 750 km)
- Ceres power spectrum deviates from the power law at λ > 270 km



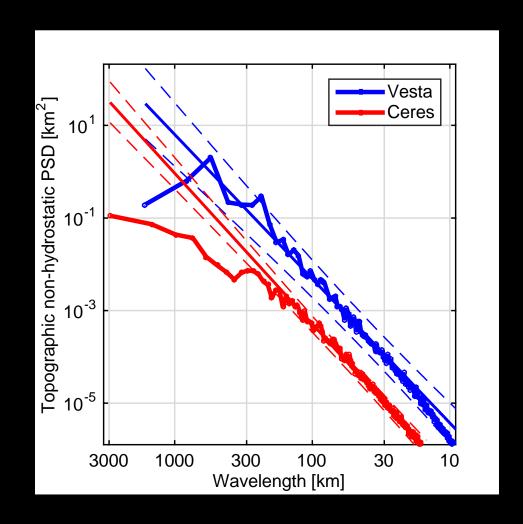


- More general approach: <u>study topography power</u> <u>spectrum</u>
- Power spectra for Vesta closely fits with the power law to the lowest degrees (λ < 750 km)
- Ceres power spectrum deviates from the power law at λ > 270 km





- More general approach: <u>study topography power</u> <u>spectrum</u>
- Power spectra for Vesta closely fits with the power law to the lowest degrees (λ < 750 km)
- Ceres power spectrum deviates from the power law at λ > 270 km





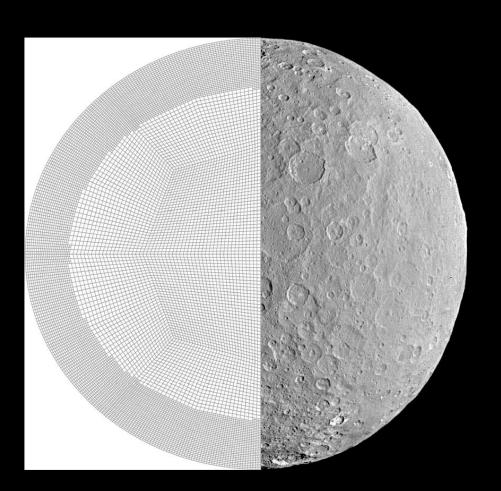








Finite element model



- Assume a density and rheology structure
- Solve Stokes equation for an incompressible flow using deal.ii library
- Compute the evolution of the outer surface power spectrum

Fu et al., 2014; Fu et al, 2017



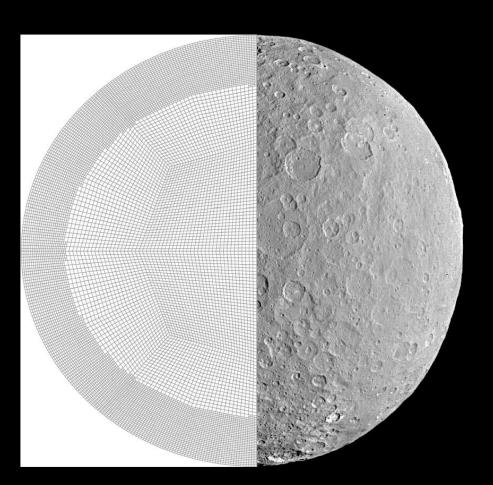








Finite element model



- Assume a density and rheology structure
- Solve Stokes equation for an incompressible flow using deal.ii library
- Compute the evolution of the outer surface power spectrum

Fu et al., 2014; Fu et al, 2017



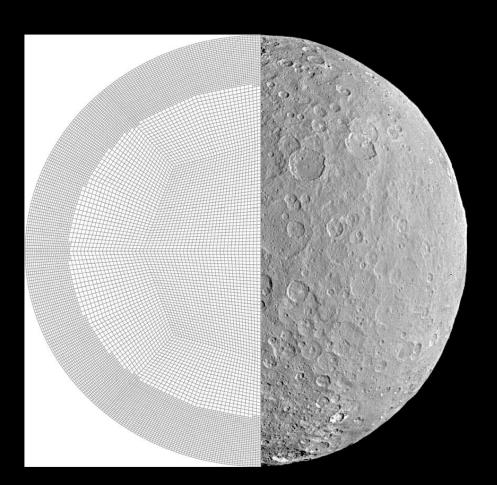








Finite element model

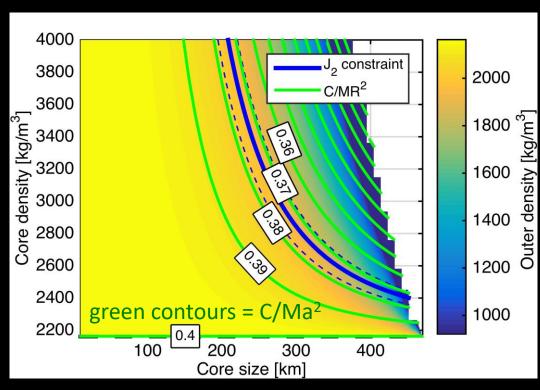


- Assume a density and rheology structure
- Solve Stokes equation for an incompressible flow using deal.ii library
- **Compute the evolution** of the outer surface power spectrum

Fu et al., 2014; Fu et al, 2017

Ceres internal structure

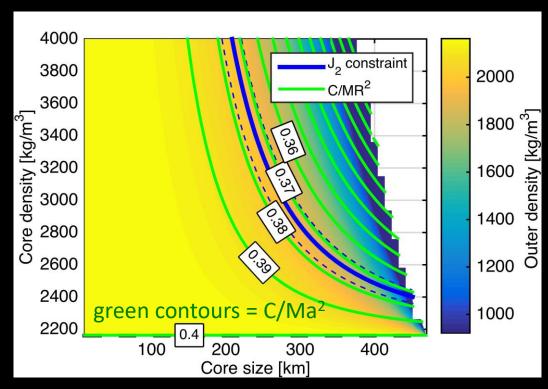
 Simplest model to interpret the gravitytopography data



Using Tricarico 2014 for computing hydrostatic equilibrium



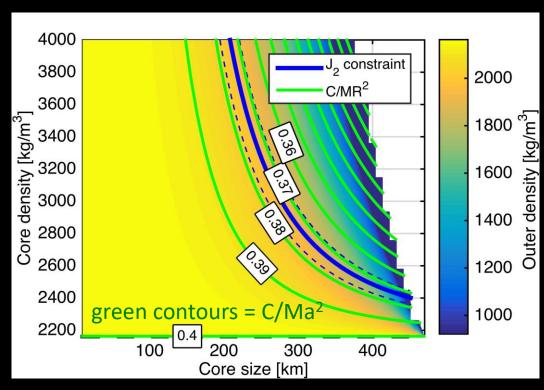
- Simplest model to interpret the gravitytopography data
- Only 5 parameters: two densities, two radii and rotation rate



Using Tricarico 2014 for computing hydrostatic equilibrium



- Simplest model to interpret the gravitytopography data
- Only 5 parameters: two densities, two radii and rotation rate
- Yields $C/Ma^2 = 0.373$ $C/M(R_{\text{yol}})^2 = 0.392$

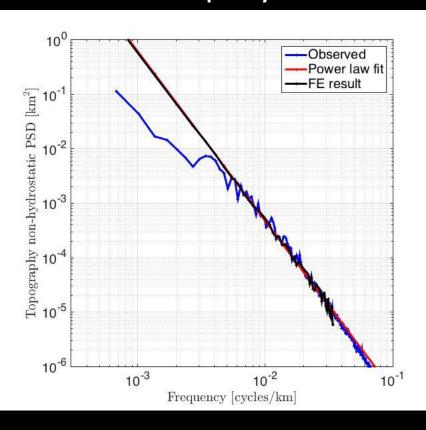


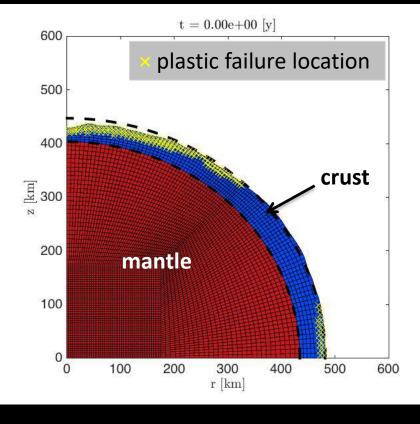
Using Tricarico 2014 for computing hydrostatic equilibrium



relaxation in the frequency domain

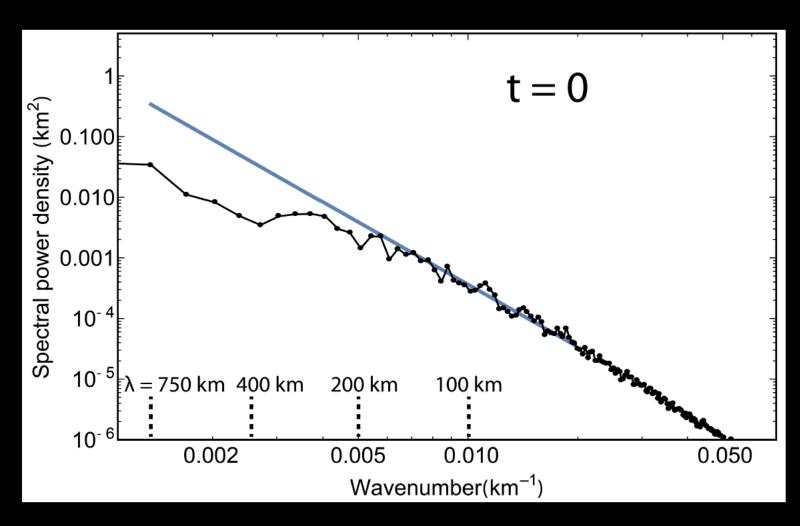
relaxation in the spatial domain





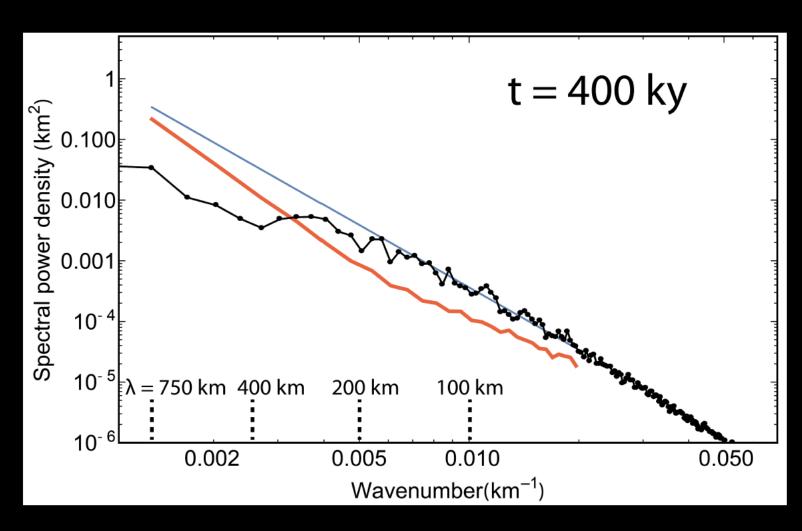






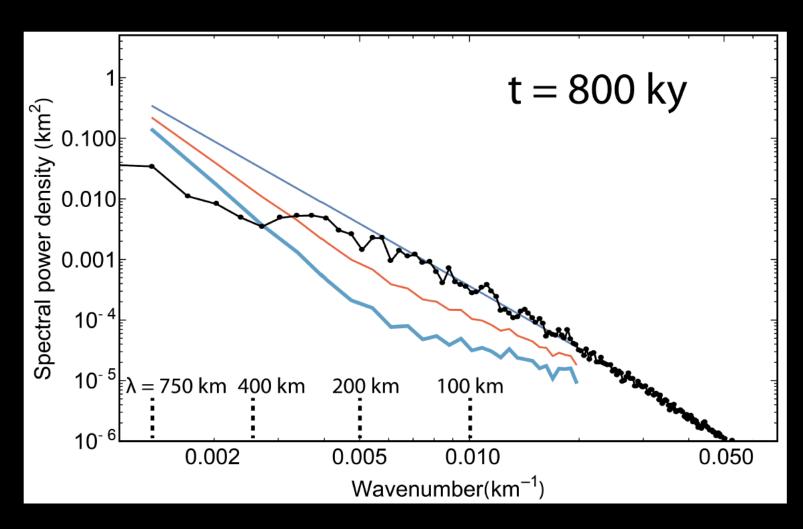


Ice shell, rocky interior



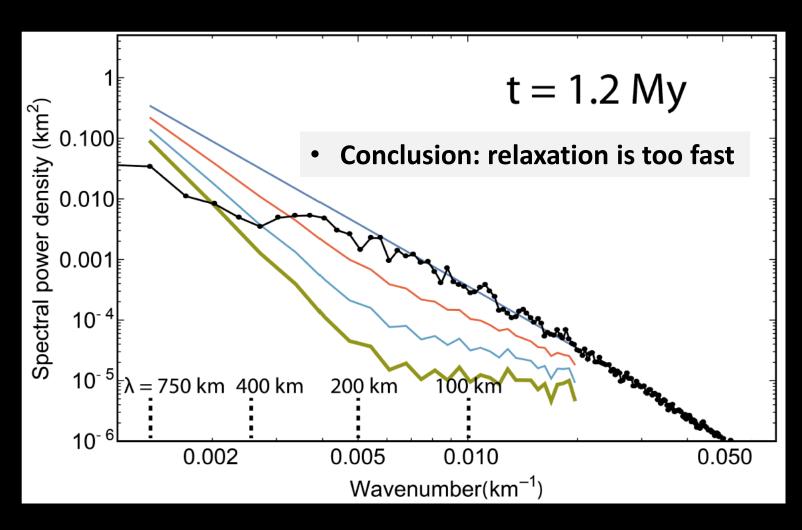


Ice shell, rocky interior



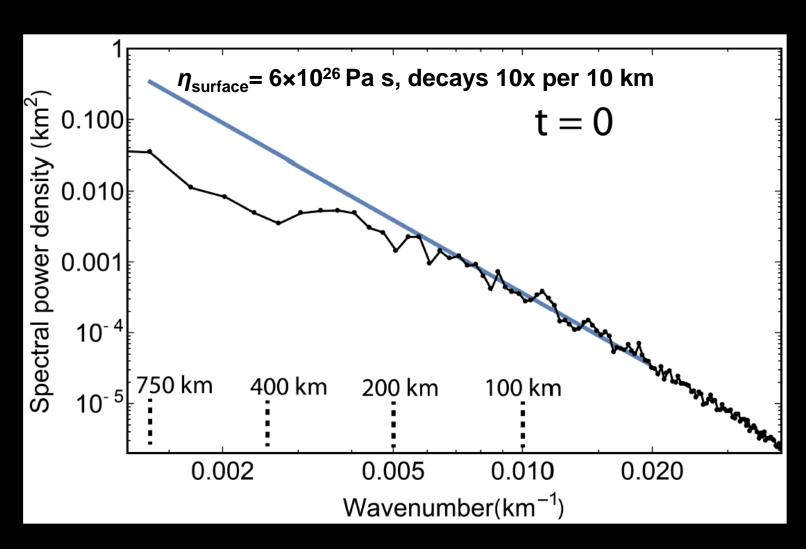


Ice shell, rocky interior

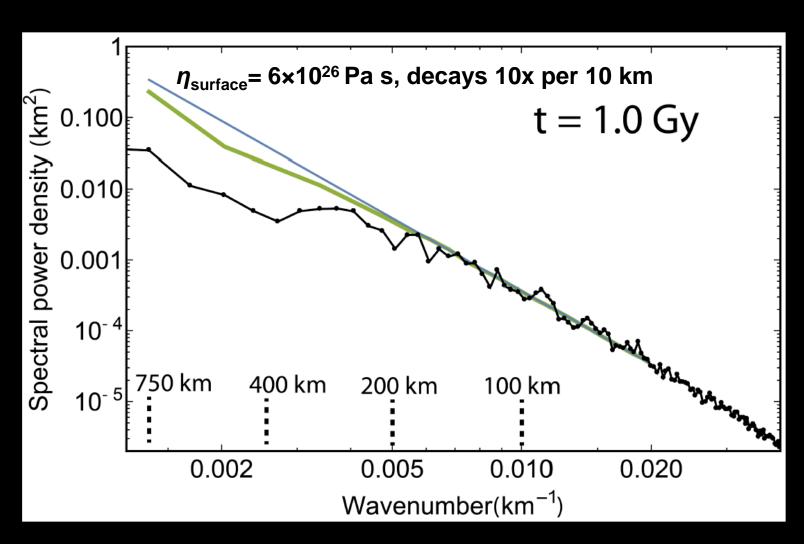




Stiff surface, weak interior

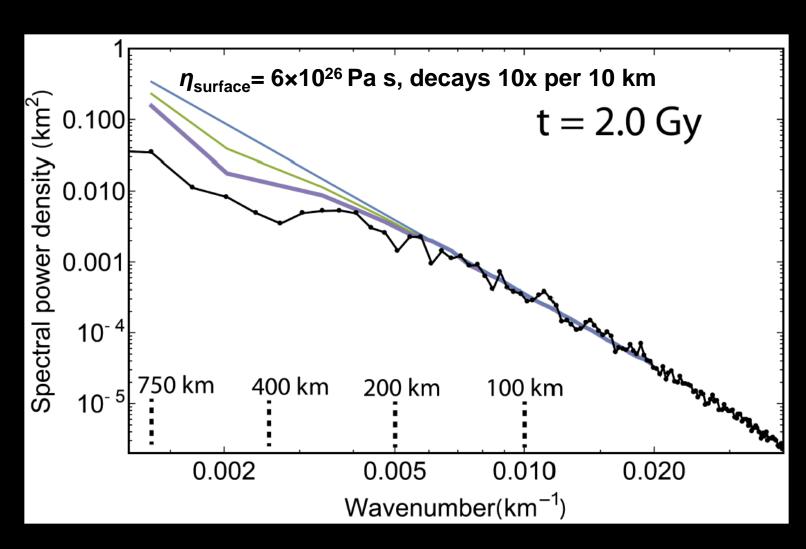




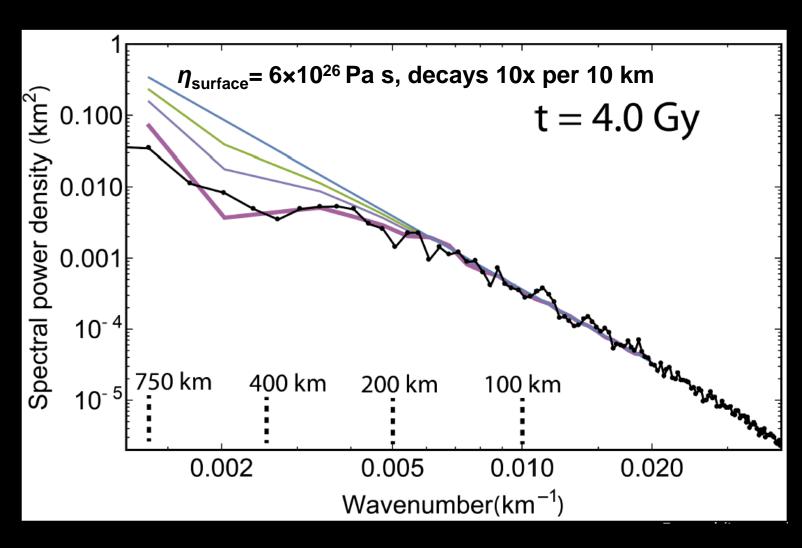




Stiff surface, weak interior





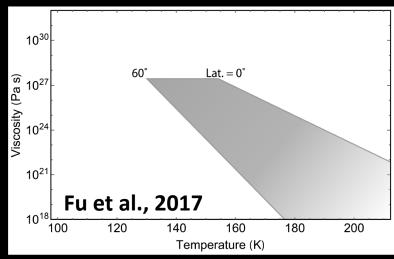


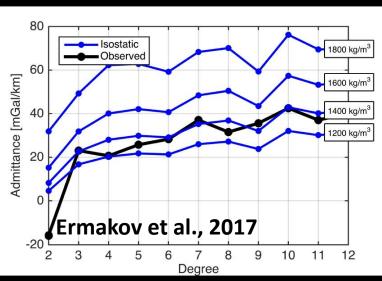


Rheology constraint from FE modeling

Density constraint from admittance modeling

Rheology and density constraints





- Ceres crust is ~ 1000 times stronger than water ice
- ➤ Must be dominated by rocklike materials. Water ice in the Ceres' crust <35 vol%
- Crust dominated salt and clathrates phases
- Low core density implies its hydrated state

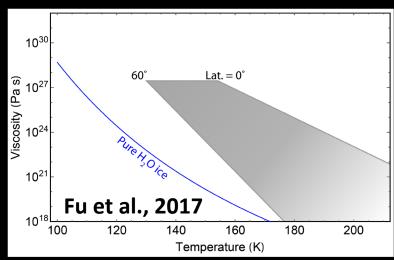


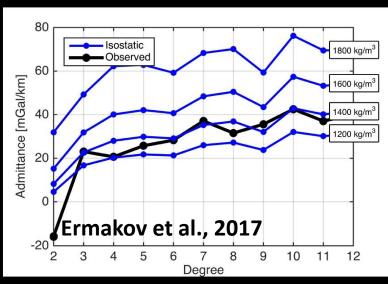
Rheology constraint modeling

Density constraint from









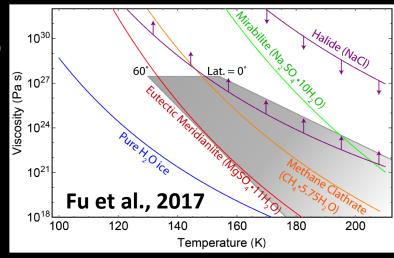
- Ceres crust is ~ 1000 times stronger than water ice
- Must be dominated by rocklike materials. Water ice in the Ceres' crust <35 vol%
- Crust dominated salt and clathrates phases
- Low core density implies its hydrated state

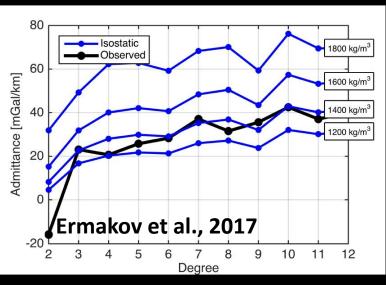


Rheology constraint from FE modeling

Density constraint from admittance modeling

Rheology and density constraints





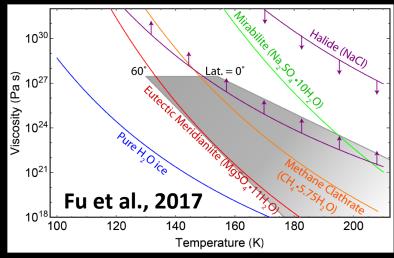
- Ceres crust is ~ 1000 times stronger than water ice
- ➤ Must be dominated by rocklike materials. Water ice in the Ceres' crust <35 vol%
- Crust dominated salt and clathrates phases
- Low core density implies its hydrated state

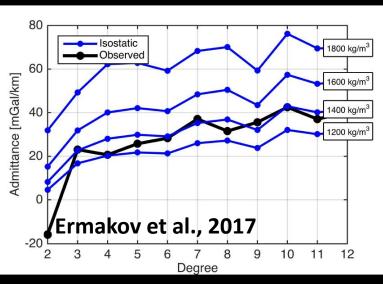


Rheology constraint from FE modeling

Density constraint from admittance modeling

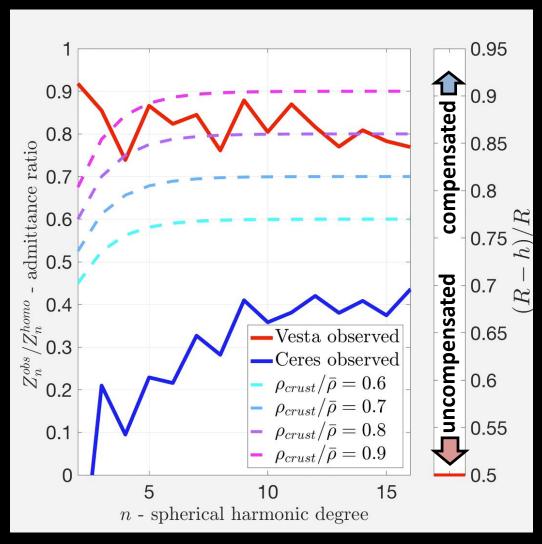
Rheology and density constraints





- Ceres crust is ~ 1000 times stronger than water ice
- ➤ Must be dominated by rocklike materials. Water ice in the Ceres' crust <35 vol%
- Crust dominated salt and clathrates phases
- Low core density implies its hydrated state











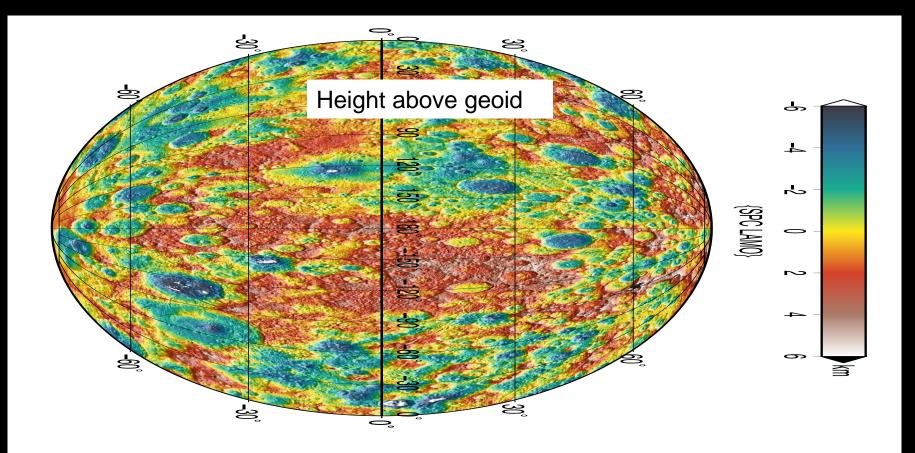




- Vesta topography is uncompensated
- Vesta acquired most of its topography when the crust was already cool and not-relaxing

- Ceres topography is compensated
- Lower viscosities (compared to Vesta) enabled <u>relaxation</u> of topography to the isostatic state





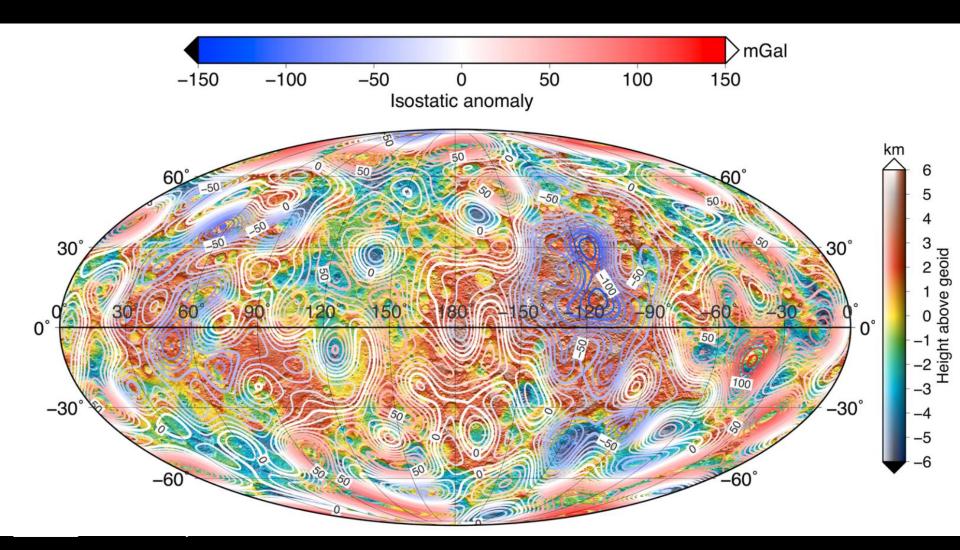
Park et al., 2016

Reference ellipsoid:

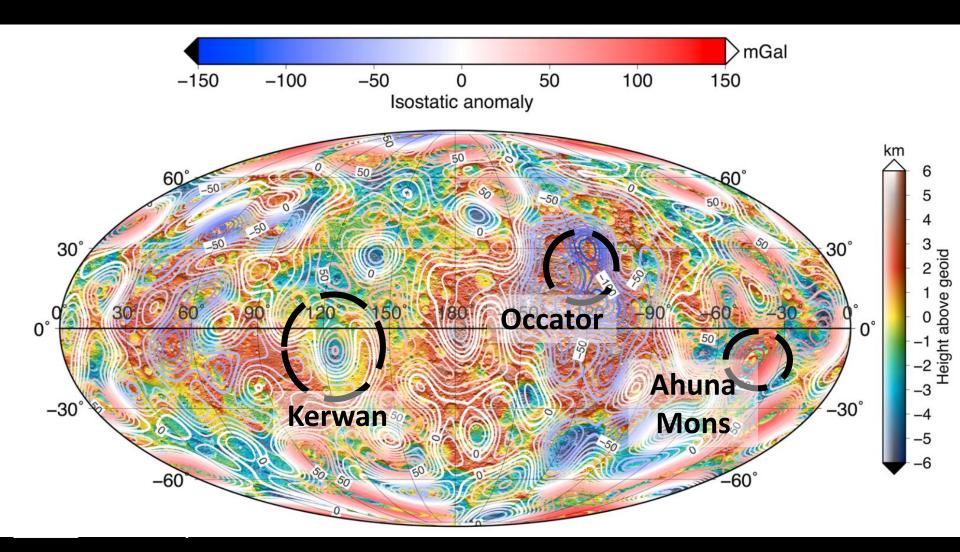
a = 445.9 km

c = 482.0 km

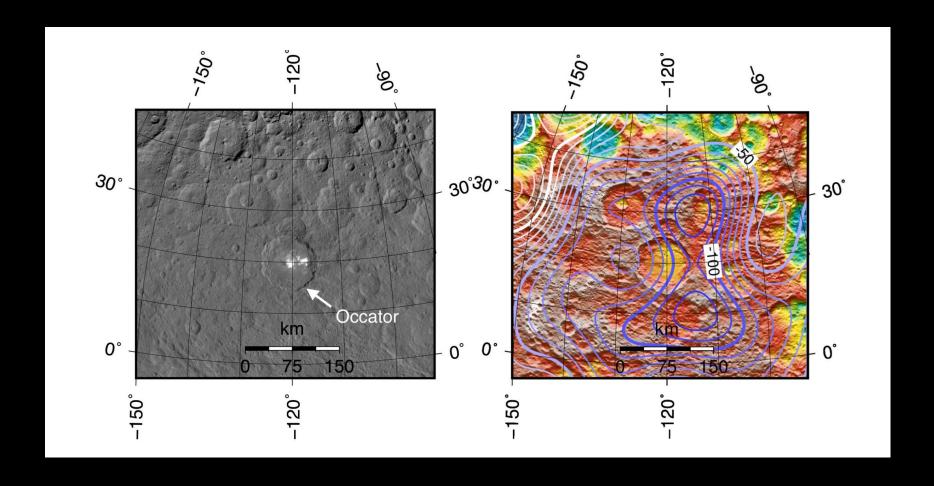






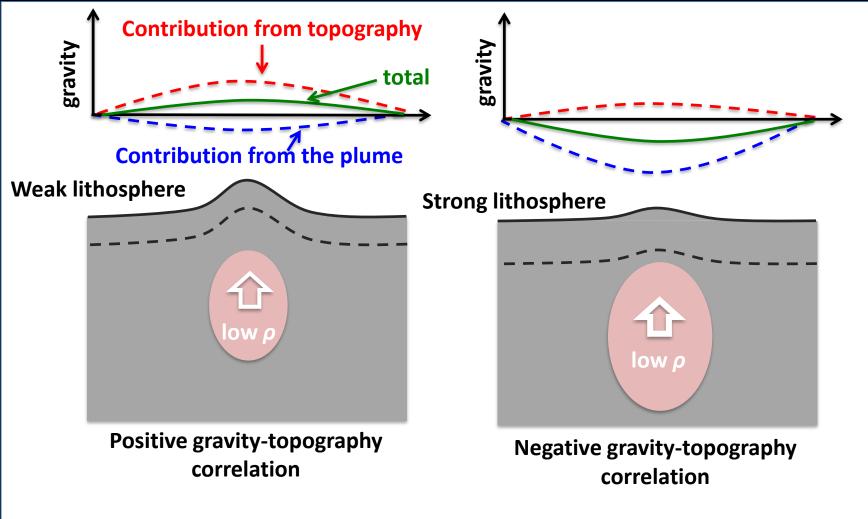






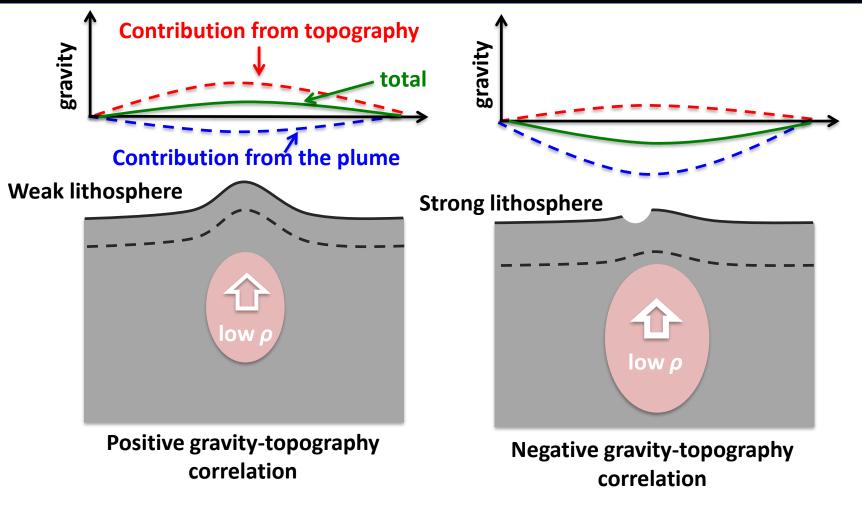






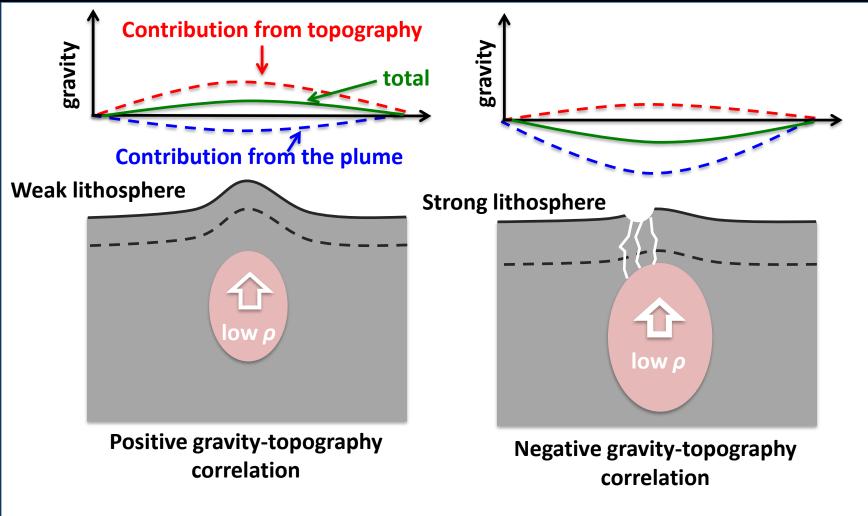










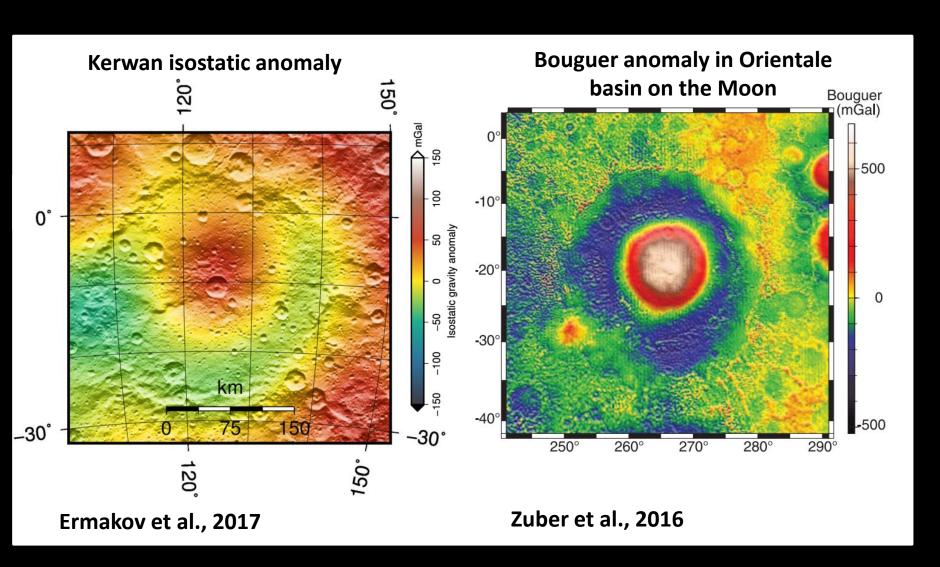




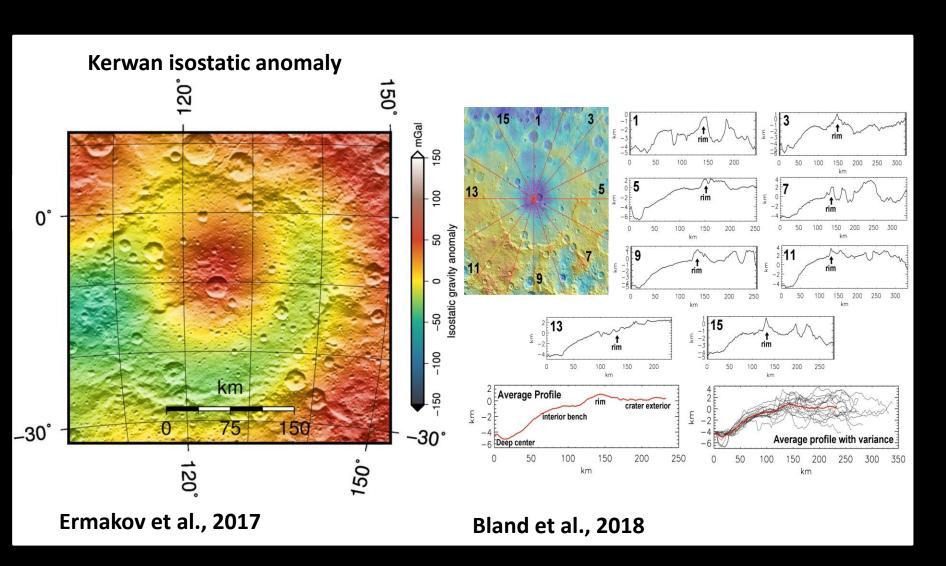












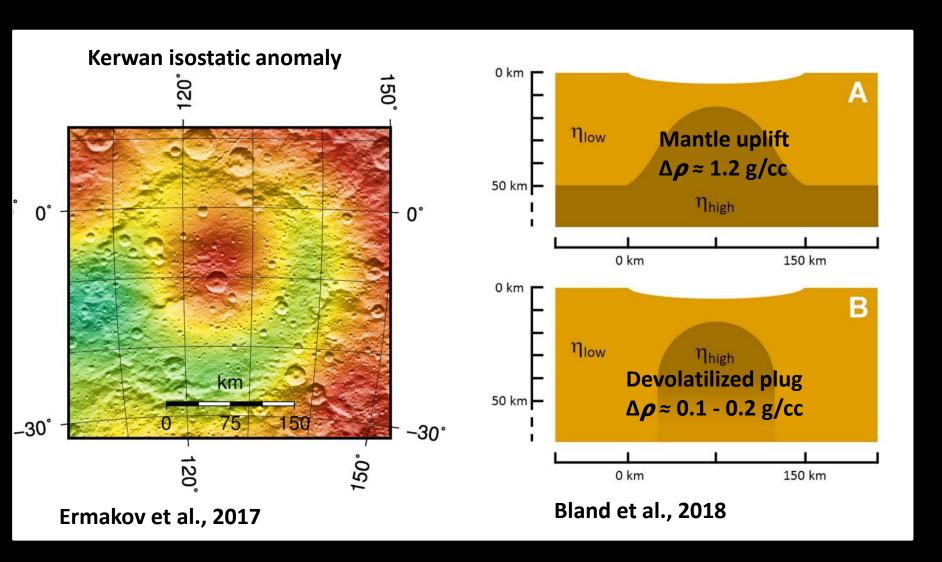


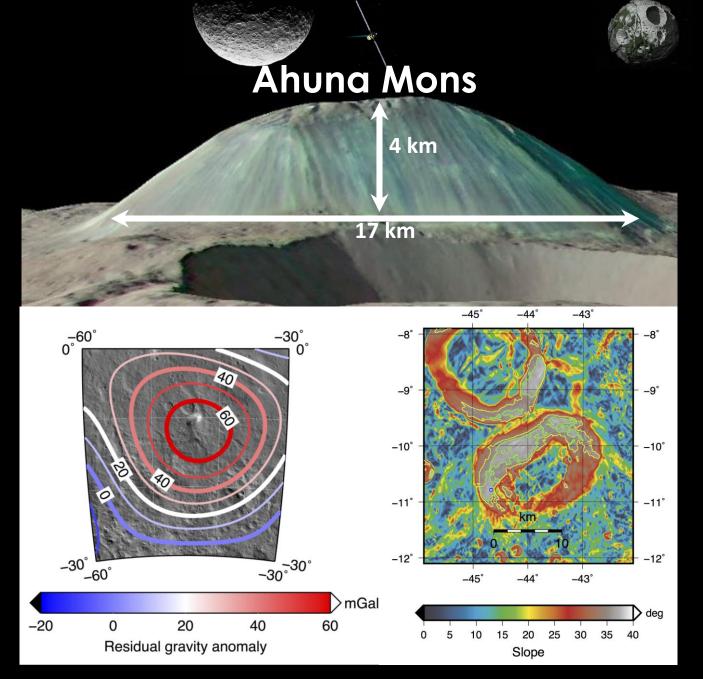






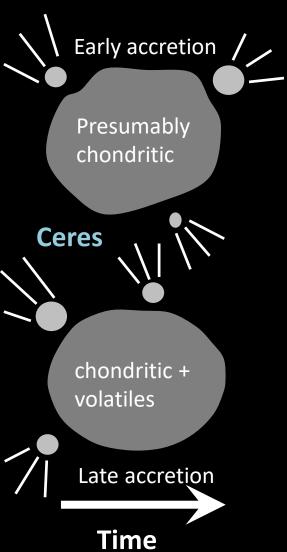












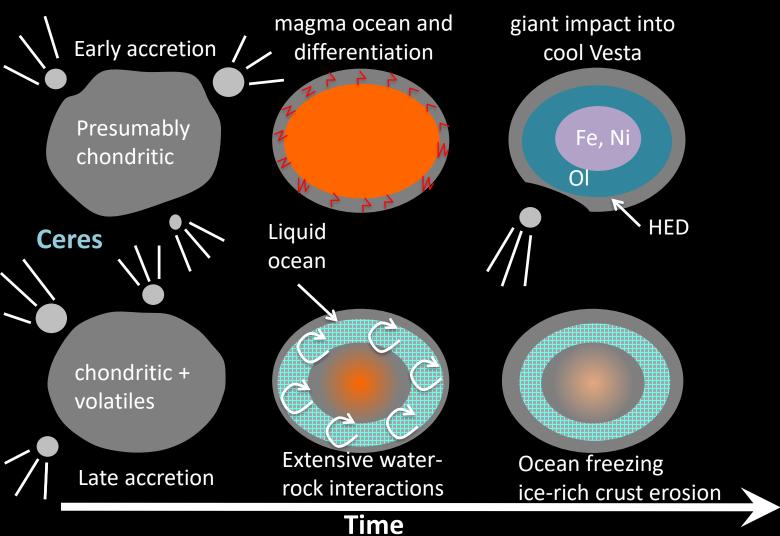


Vesta magma ocean and Early accretion differentiation Presumably chondritic Liquid Ceres ocean chondritic + volatiles Extensive water-Late accretion rock interactions \

Time

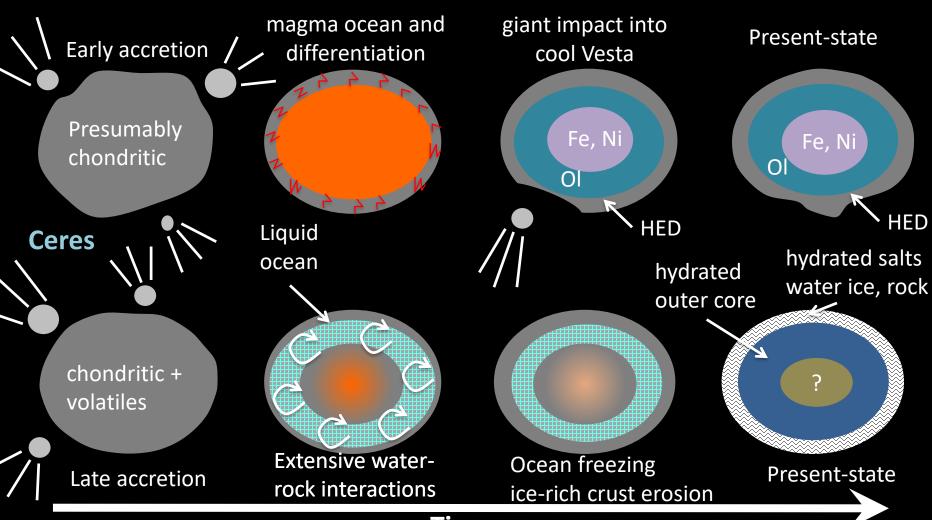
Vesta and Ceres comparative evolution

Vesta



Vesta and Ceres comparative evolution

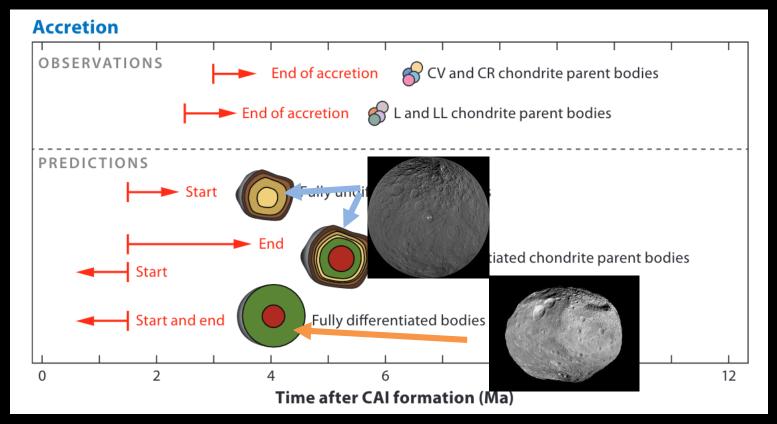




Time



A spectrum of planetesimal differentiation



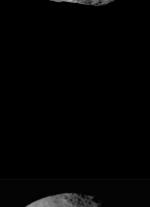
Weiss & Elkins-Tanton, 2013

Summary

- Formed early (< 5 My after CAI)
- Once hot and hydrostatic, Vesta is no longer either
- Differentiated interior
- Most of topography acquired when Vesta was already cool => uncompensated topography
- Combination of gravity/topography data with meteoritic geochemistry data provides constraints on the internal structure



- late formation
- and/or heat transfer due to hydrothermal circulation
- Partially differentiated interior
- Experienced viscous relaxation
- Much lower surface viscosities (compared to Vesta) allowed compensated topography
- Ceres' crust is light (based on admittance analysis) and strong (based on FE relaxation modeling)
- Not much water ice in Ceres crust (<35 vol%) now



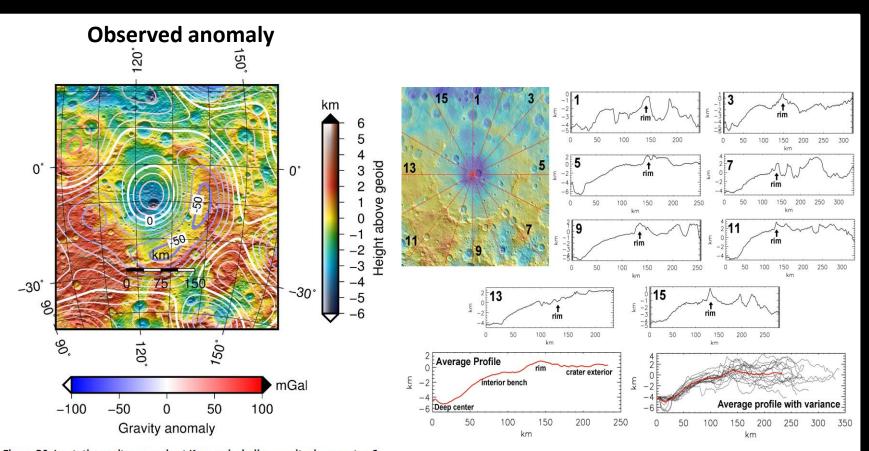


Figure S6. Isostatic gravity anomaly at <u>Kerwan</u> including gravity degree 3 to 16. Contours are for gravity anomaly and the contour interval is 10 mGal. Colors show relative topography in the <u>Kerwan</u> region. After Ermakov et al. 2017.

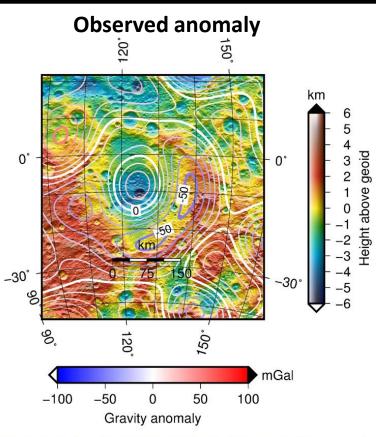


Figure S6. Isostatic gravity anomaly at <u>Kerwan</u> including gravity degree 3 to 16. Contours are for gravity anomaly and the contour interval is 10 mGal. Colors show relative topography in the <u>Kerwan</u> region. After Ermakov et al. 2017.

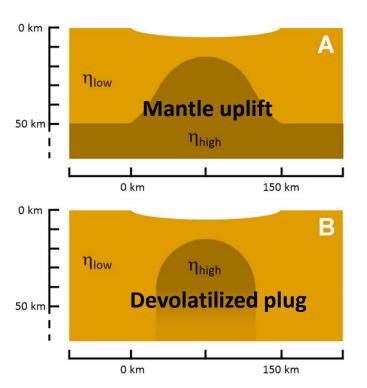
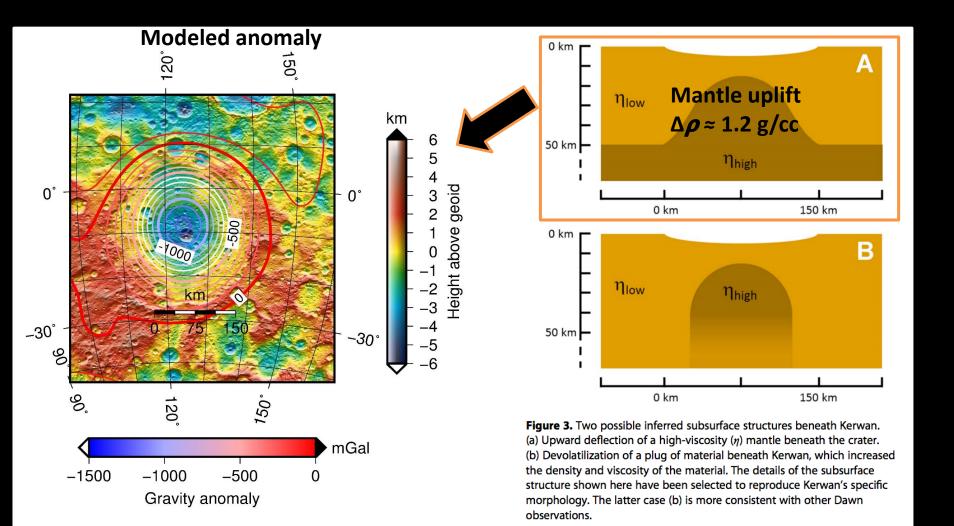
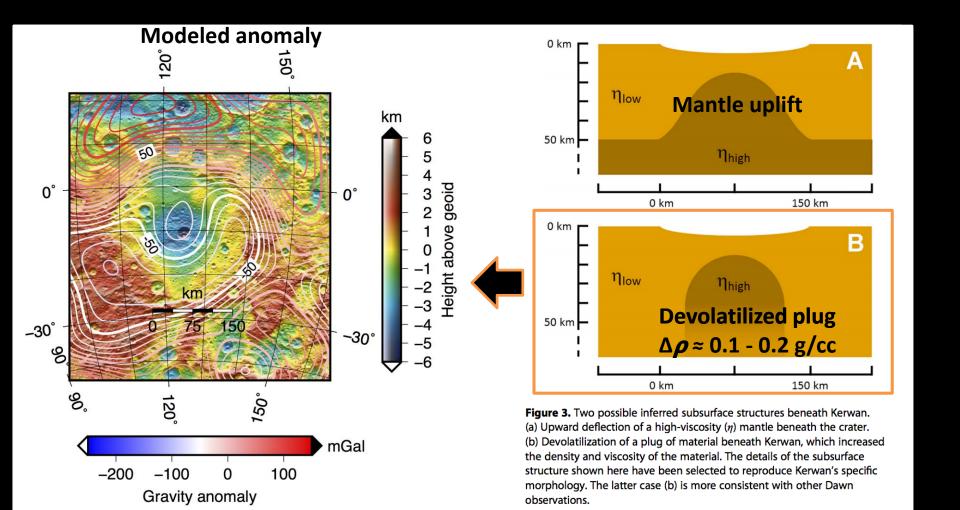
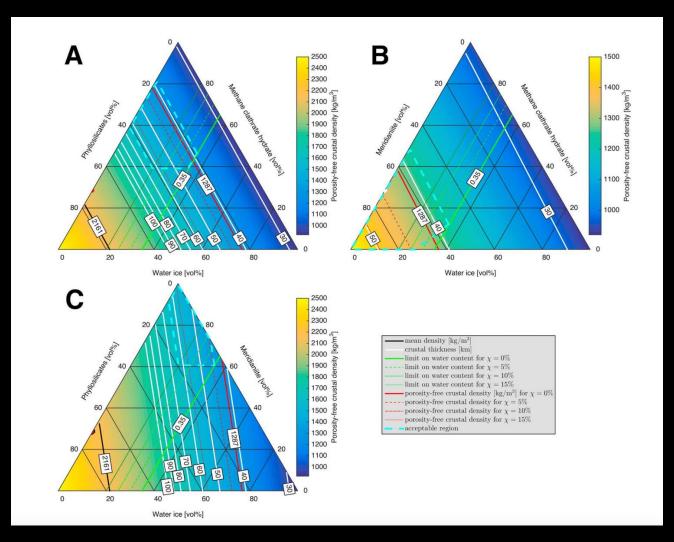


Figure 3. Two possible inferred subsurface structures beneath Kerwan. (a) Upward deflection of a high-viscosity (η) mantle beneath the crater. (b) Devolatilization of a plug of material beneath Kerwan, which increased the density and viscosity of the material. The details of the subsurface structure shown here have been selected to reproduce Kerwan's specific morphology. The latter case (b) is more consistent with other Dawn observations.





Crustal composition constraints



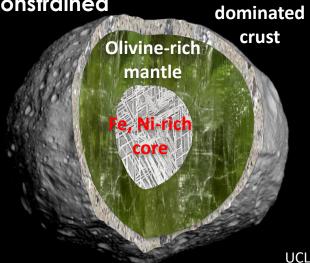
Ermakov et al., 2017

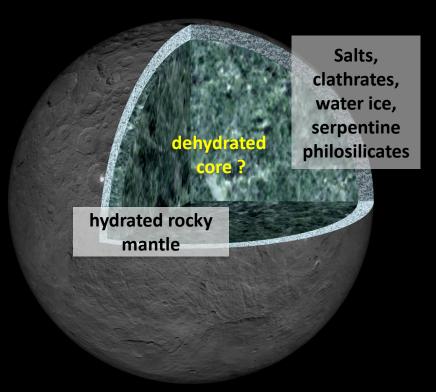
Internal structures of Vesta and Ceres

Ceres→

- Crust is light (1.1-1.4 g/cc) and mechanically rocklike w
- Mantle density ~2.4 g/cc and unlithified at least to a depth of 100 km

Possible dehydrated rocky core remains HED-unconstrained



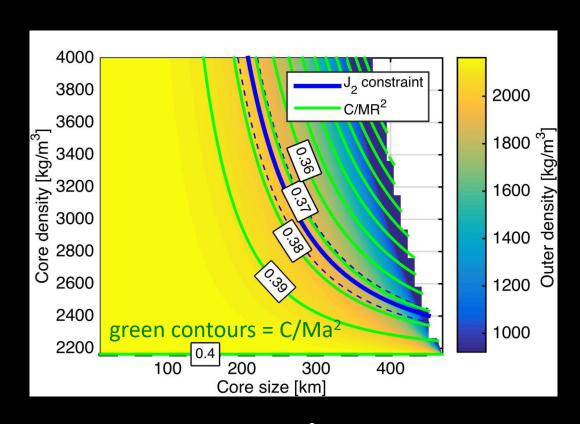


←Vesta

- Crustal density constrained by HEDs and admittance (2.8 g/cc)
- Assuming density of iron meteorites (5-8 g/cc), the core radius is 110 155 km

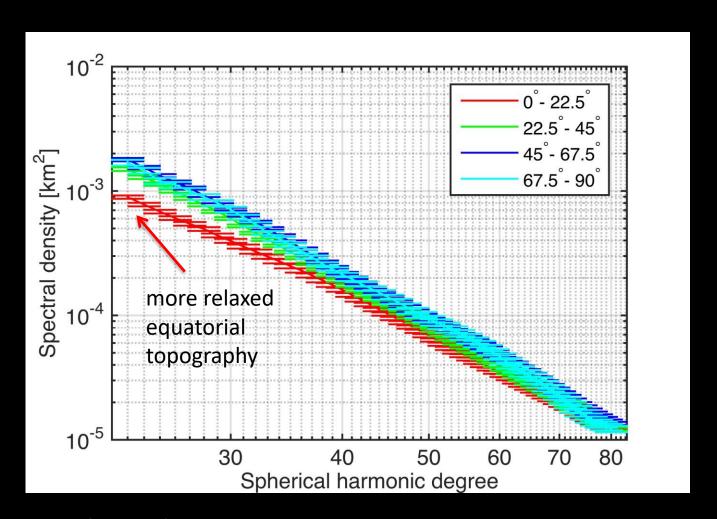
Two-layer model

- Simplest model to interpret the gravitytopography data
- Only 5 parameters: two densities, two radii and rotation rate
- Yields $C/Ma^2 = 0.373$ $C/M(R_{vol})^2 = 0.392$



Using Tricarico 2014 for computing hydrostatic equilibrium

Latitude dependence of relaxation



Ermakov et al., in prep

Gravity and topography in spherical harmonics

Shape radius vector

Gravitational potential

Power Spectral Density

$$S_n^{gg} = \mathop{\mathring{o}}_{m=0}^n \frac{C_{nm}^2 + S_{nm}^2}{2n+1}$$

gravity

$$S_{n}^{tt} = \mathop{a}\limits_{m=0}^{n} \frac{A_{nm}^{2} + B_{nm}^{2}}{2n+1}$$

topography

$$S_n^{gg} = \mathop{\mathring{o}}_{m=0}^n \frac{C_{nm}^2 + S_{nm}^2}{2n+1}$$

$$S_n^{tt} = \mathop{\mathring{o}}_{m=0}^n \frac{A_{nm}^2 + B_{nm}^2}{2n+1}$$

$$S_n^{gt} = \mathop{\mathring{o}}_{m=0}^n \frac{A_{nm}C_{nm} + B_{nm}S_{nm}}{2n+1}$$

gravity-topography cross power

Isostatic model

Z_n - gravity-topography admittance

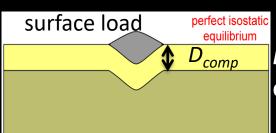
$$Z_n = \frac{S_{gt}}{S_{tt}}$$

Linear two-layer hydrostatic model

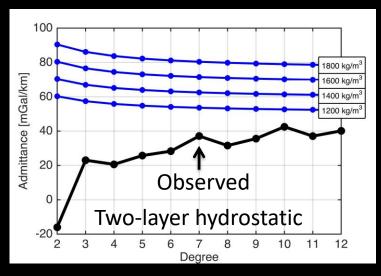
$$Z_n = \frac{GM}{R^3} \frac{3(n+1)}{2n+1} \frac{\Gamma_{crust}}{\Gamma_{mean}}$$

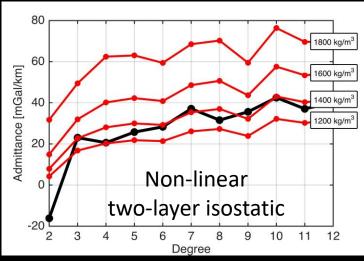
> Linear isostatic model

$$Z_{n} = \frac{GM}{R^{3}} \frac{3(n+1)}{2n+1} \frac{\Gamma_{crust}}{\Gamma_{mean} \hat{\mathfrak{g}}} \hat{\mathfrak{g}}^{1} - \hat{\mathfrak{g}}^{1} - \frac{D_{comp}}{R} \hat{\mathfrak{g}}^{0} \hat{\mathfrak{g}}^{1}$$

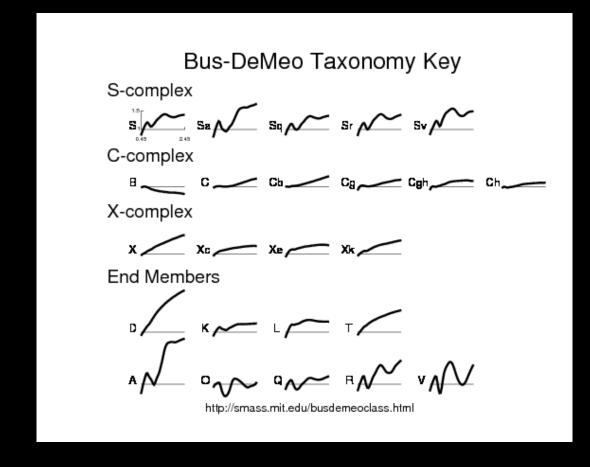


D_{comp}- depth of compensation

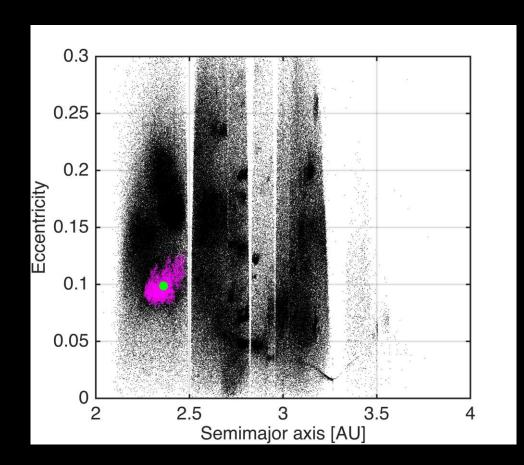




Unique basaltic spectrum



- Unique basaltic spectrum
- A group of asteroids in the dynamical vicinity of Vesta with similar spectra



- Unique basaltic spectrum
- A group of asteroids in the dynamical vicinity of Vesta with similar spectra
- Large depression in the southern hemisphere of Vesta

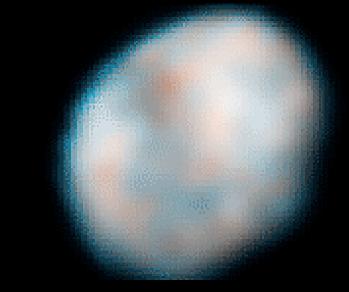
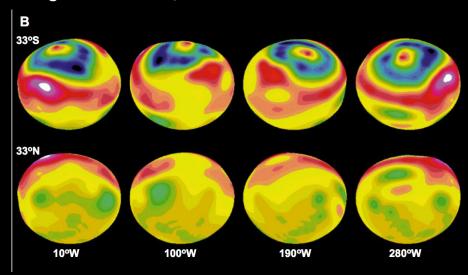
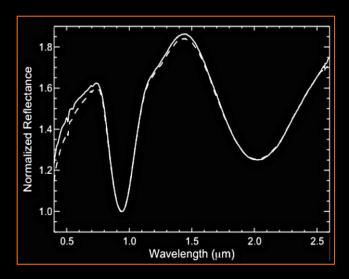


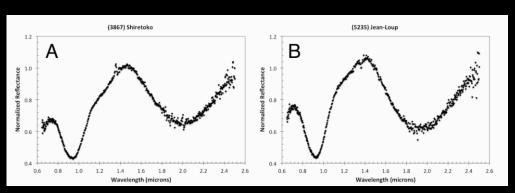
Image credit: NASA/HST



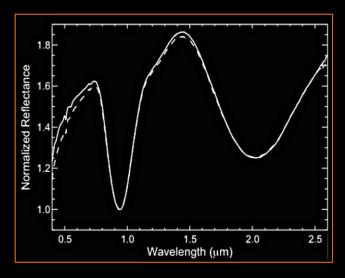
- Unique basaltic spectrum
- A group of asteroids in the dynamical vicinity of Vesta with similar spectra
- Large depression in the southern hemisphere of Vesta
- A group of Howardite-Eucrite-Diogenite (HED) meteorites, with similar reflectance spectra



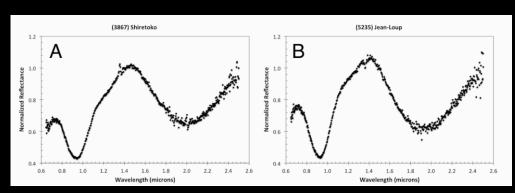
- ↑ Reflectance spectra of eucrite Millbillillie from Wasson et al. (1998)
- **V**-type asteroids spectra from Hardensen et al., (2014)



- Unique basaltic spectrum
- A group of asteroids in the dynamical vicinity of Vesta with similar spectra
- Large depression in the southern hemisphere of Vesta
- A group of Howardite-Eucrite-Diogenite (HED) meteorites, with similar reflectance spectra
- Strongest connection between a class of meteorites and an asteroidal family



- ↑ Reflectance spectra of eucrite Millbillillie from Wasson et al. (1998)
- **V**-type asteroids spectra from Hardensen et al., (2014)



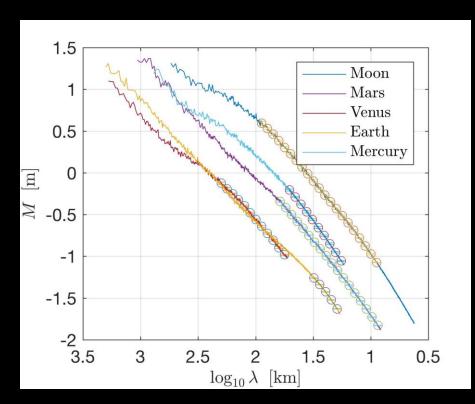
Note on Vening-Meinesz and Kaula rules

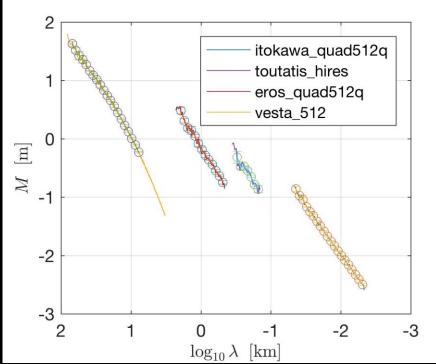
- Vening-Meinesz rule for variance of topography (Vening-Meinesz, 1951)
 V_t ~ 1/n²
- Kaula law for RMS of gravity (Kaula, 1963) $M_g \sim 1/n^2$
- Are these two rules consistent assuming uncompensated topography?

$$V_t \sim 1/n^2 => M_t \sim 1/n^{1.5} => M_g \sim 1/n^{2.5}$$

- But Kaula rule says M_g ~ 1/n² NOT M_g ~ 1/n^{2.5}
- Typically assumed in the literature Kaula and Vening-Meinesz rules are not mutually consistent assuming uncompensated topography

RMS spectra





Power laws

General form of a power law

$$M=AR^{\alpha_1}\varrho^{\alpha_2}\lambda^{\alpha_3}$$

• Power law assuming (inverse) surface gravity scaling $(g \sim R^* \rho)$

$$M=AR^{-1}\varrho^{-1}\lambda^{\alpha_3}$$

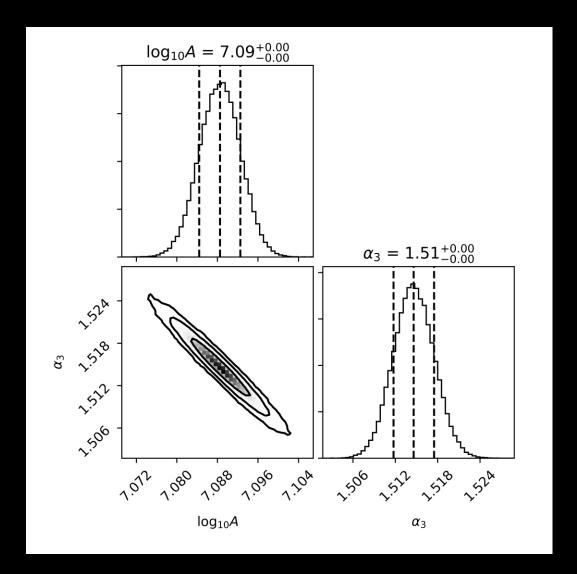
If we take a log₁₀ of M, we get an equation of a hyperplane

$$\log_{10}M = \log_{10}A + \alpha_1\log_{10}R + \alpha_2\log_{10}Q + \alpha_3\log_{10}\lambda$$

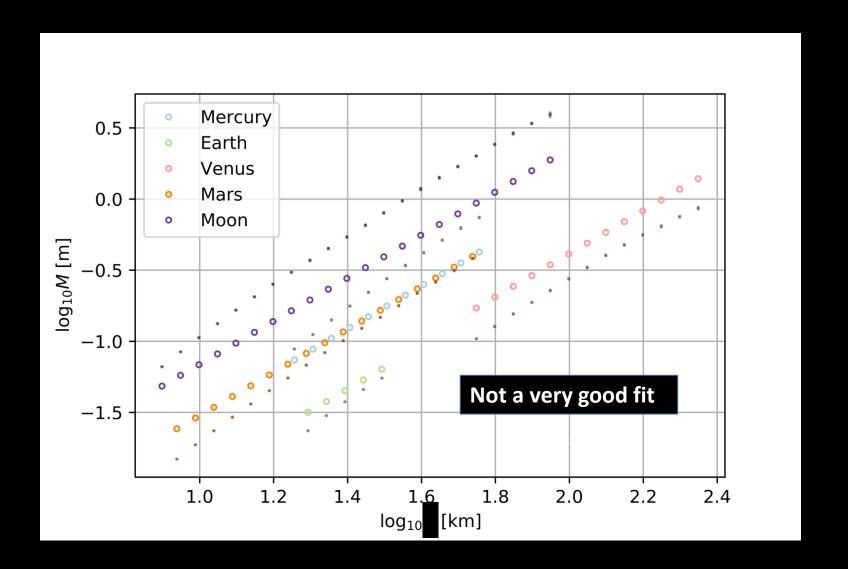
- In our data set, we have a lot of points along the λ direction and not as many points on the other two (R and ρ) directions.
 - In the R and ho directions, we have as many data points as we have bodies
 - In the λ direction, we have as many data points as many we have λ bins.

Results of the MCMC runs

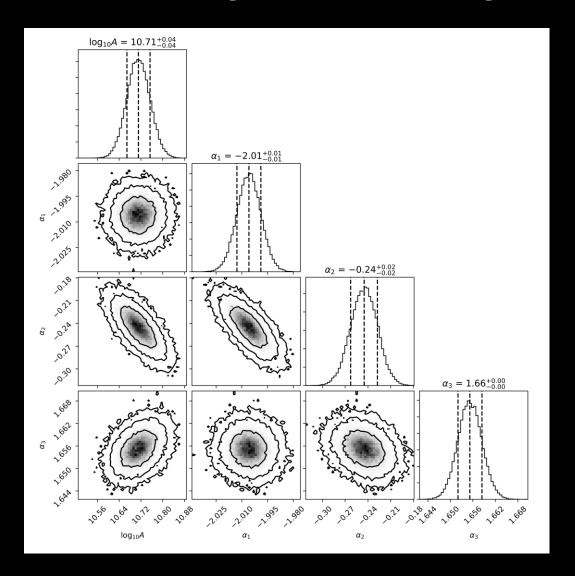
Planets, gravity scaling



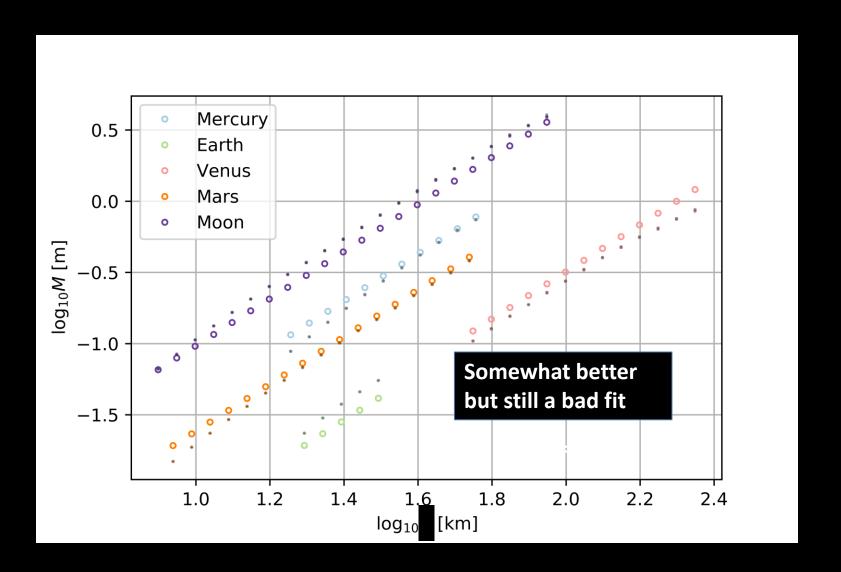
Planets, gravity scaling



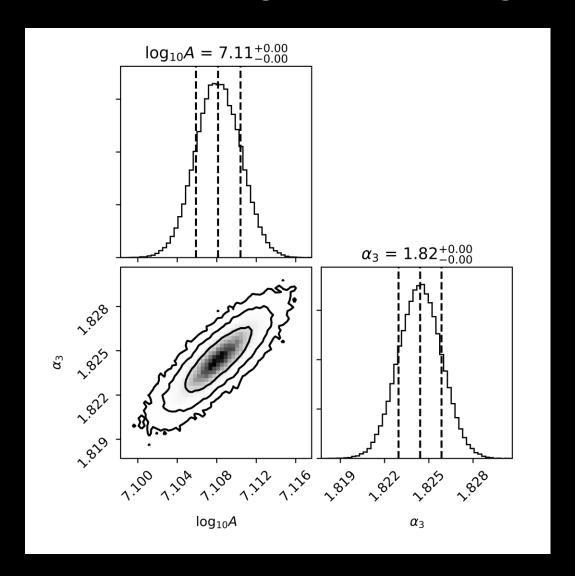
Planets, general scaling



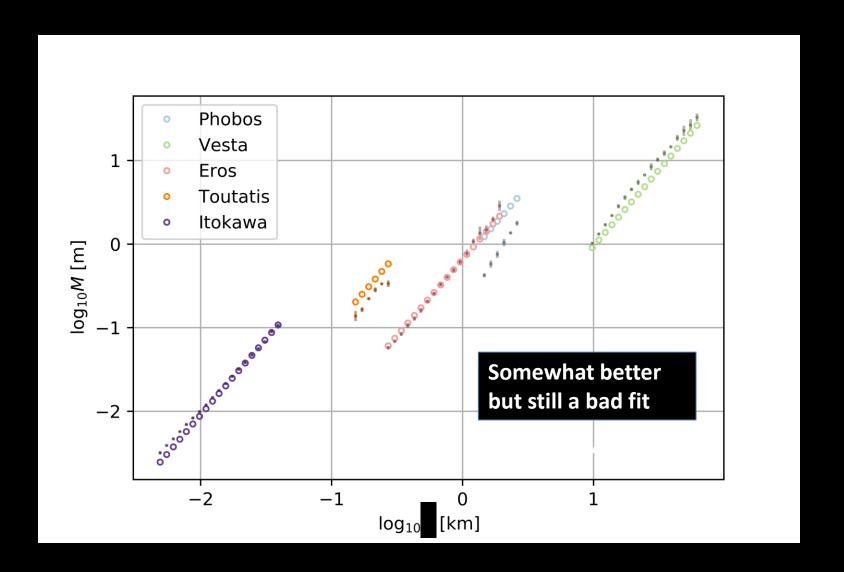
Planets, general scaling



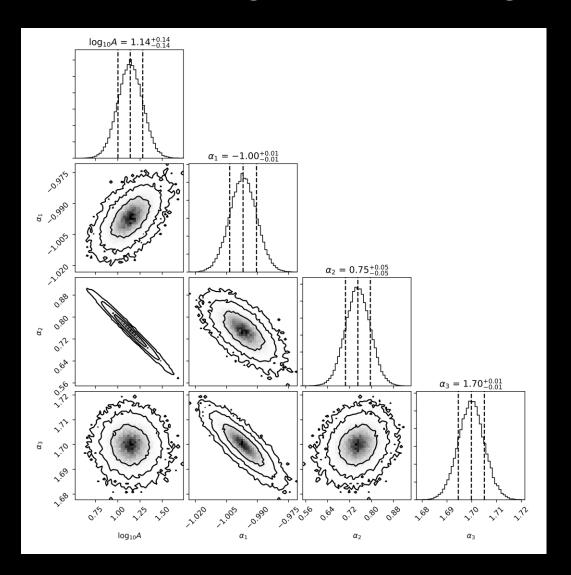
Asteroids, gravity scaling



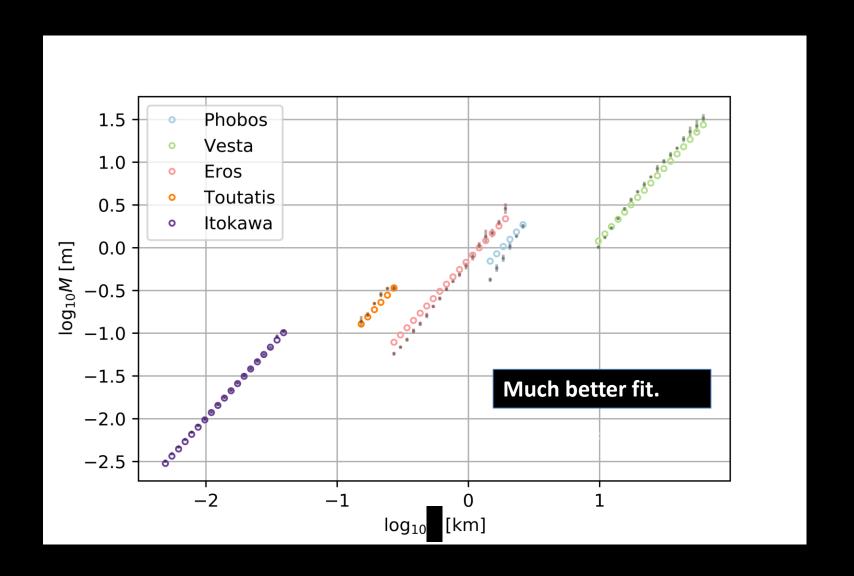
Asteroids, gravity scaling



Asteroids, general scaling



Asteroids, general scaling



A priori constraint on gravity RMS

Choose R and ρ

Given R and ρ and a range of λ , sample multivariate normal distribution to get A, $\alpha_1, \alpha_2, \alpha_3$

Find the upper and lower bounds on the gravity RMS spectum

Given A, α_1 , α_2 , α_3 , compute topography RMS spectrum

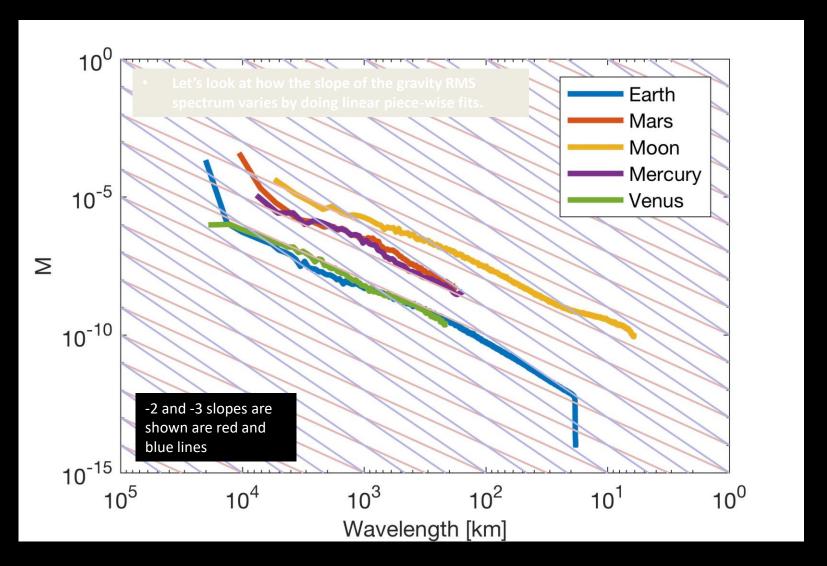
Given topography RMS spectrum, generate SH coefficients that follow the chosen spectrum

Compute gravity-fromtopography using Wieczorek & Phillips 1998 until convergence w.r.t. to the power of topography

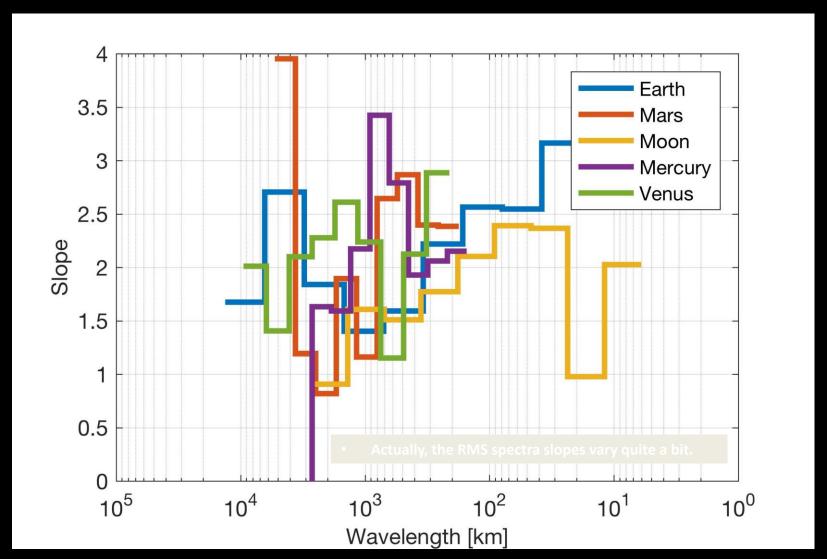
Summary

- Topography RMS spectra of 4 terrestrial planets and the Moon cannot be simultaneously fit with a single power law of the gravity-scaling or general form.
- Topography RMS spectra of asteroids CANNOT be satisfactorily fit with a power law the gravity-scaling form.
- Topography RMS spectra of asteroids CAN be satisfactorily fit with a power law of the general form.
- Despite having different internal structure, composition and mechanical properties of the surface layer, the asteroid topography spectra can be effectively modeled as a general power law

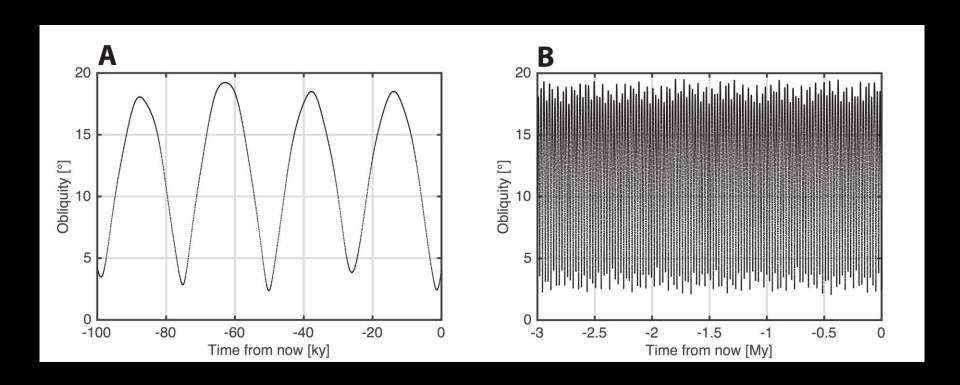
Gravity RMS spectra



Slopes of piecewise fitted gravity RMS spectra

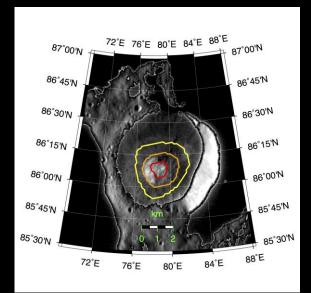


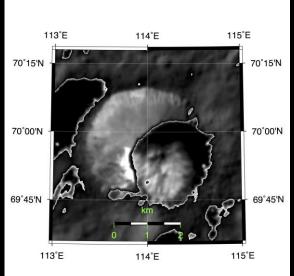
Ceres' obliquity history

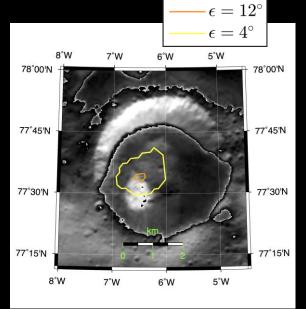


- Obliquity varies between 2.4° and 19.7°
- The main period is 24.5 ky
- We happen to visit Ceres when its obliquity is minimal

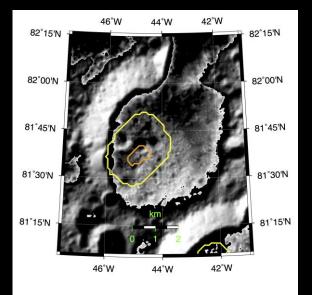
Bright Crater Floor Deposits (BCFDs)

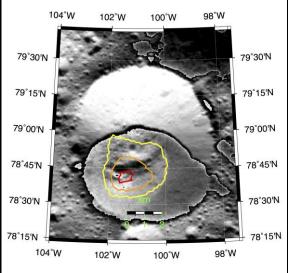


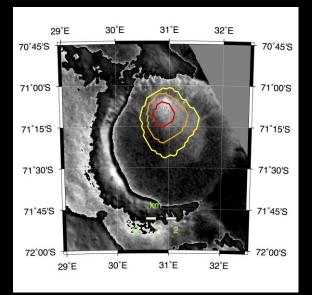




 $\epsilon=20^{\circ}$





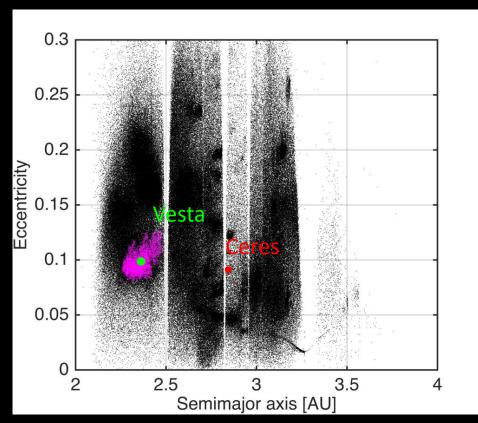


13

Why Ceres?

- Largest body in the asteroid belt
- Low density implies high volatile content
- Conditions for subsurface ocean
- Much easier to reach than other ocean worlds

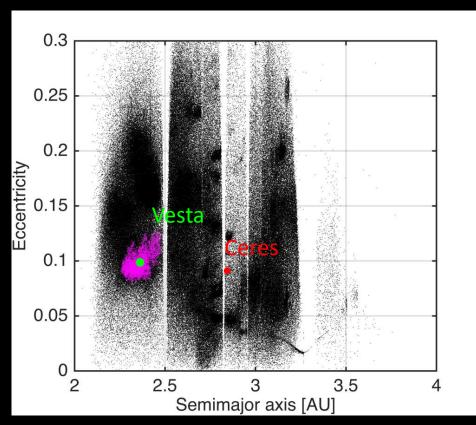
Ceres location in the asteroid belt



Why Ceres?

- Largest body in the asteroid belt
- Low density implies high volatile content
- Conditions for subsurface ocean
- Much easier to reach than other ocean worlds
- Major unexplored object in the asteroid belt

Ceres location in the asteroid belt



What did we know before Dawn

Castillo-Rogez and McCord 2010

Ceres accreted as a mixture of ice and rock just a few My after the condensation of Calcium Aluminum-rich Inclusions (CAIs), and later differentiated into a water mantle and a mostly anhydrous silicate core.

What did we know before Dawn

Castillo-Rogez and McCord 2010

Ceres accreted as a mixture of ice and rock just a few My after the condensation of Calcium Aluminum-rich Inclusions (CAIs), and later differentiated into a water mantle and a mostly anhydrous silicate core.

Zolotov 2009

Ceres formed relatively late from planetesimals consisting of hydrated silicates.

What did we know before Dawn

Castillo-Rogez and McCord 2010

Ceres accreted as a mixture of ice and rock just a few My after the condensation of Calcium Aluminum-rich Inclusions (CAIs), and later differentiated into a water mantle and a mostly anhydrous silicate core.

Zolotov 2009

Ceres formed relatively late from planetesimals consisting of hydrated silicates.

Bland 2013

If Ceres does contain a water ice layer, its warm diurnallyaveraged surface temperature ensures extensive viscous relaxation of even small impact craters especially near equator